

3D Disruption Mitigation Modeling with M3D-C1

by

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Modeling of Disruption Dynamics and Mitigation Requires a Multiphysics Model

- **Disruptions pose a risk of damage to future tokamaks, necessitating robust mitigation techniques**
- **Most promising method uses pellet injection of impurities to radiate stored energy**
- **Simulations, validated against mitigation experiments, are required to project techniques to future devices**
- **Integrated model is required to capture all relevant physics**
 - Magnetohydrodynamics (MHD) for macroscopic evolution of disruption dynamics
 - Atomic physics for ionization and radiation from injected impurities
 - Drift-kinetics for phase-space evolution of runaway electron population

KPRAD* Provides Needed Atomic Physics Information

- **KPRAD solves for impurity-plasma interaction in low-density, coronal model**
 - N.B. *not* coronal equilibrium
 - Based on ADPAK rate coefficients
 - Impurity charge states and electron density evolve according to ionization and recombination

$$\frac{\partial n_z}{\partial t} + \nabla \cdot (n_z \mathbf{v}) = \nabla \cdot (D \nabla n_z) + \mathcal{I}_{z-1} n_{z-1} - (\mathcal{I}_z + \mathcal{R}_z) n_z + \mathcal{R}_{z+1} n_{z+1} + \mathcal{S}_z$$

- Thermal energy lost from plasma due to
 - Ionization
 - Line radiation
 - Bremsstrahlung radiation
 - Recombination radiation
- **Can be subcycled much faster than typical MHD time steps**

*D.G. Whyte, et al., Proc. of the 24th Euro. Conf. on Controlled Fusion and Plasma Physics, Berchtesgaden, Germany, 1997, Vol. 21A, p. 1137.

KPRAD Couples* to the M3D-C1 Temperature Equation(s)

1) Electron and all-ions temperature equations

- Dilution cooling of ions and electrons
- Electrons lose energy to ionization and radiation
- Main ions cool on electrons

$$n_e \left[\frac{\partial T_e}{\partial t} + \mathbf{v} \cdot \nabla T_e + (\Gamma - 1) T_e \nabla \cdot \mathbf{v} \right] + \sigma_e T_e = (\Gamma - 1) \left[\eta J^2 - \nabla \cdot \mathbf{q}_e - \mathcal{P}_{rad} + Q_{ei} - \Pi_e : \nabla \mathbf{v} \right]$$

$$n_* \left[\frac{\partial T_i}{\partial t} + \mathbf{v} \cdot \nabla T_i + (\Gamma - 1) T_i \nabla \cdot \mathbf{v} \right] + \sigma_* T_i = (\Gamma - 1) \left[-\nabla \cdot \mathbf{q}_* - Q_{ei} - \Pi_* : \nabla \mathbf{v} + \frac{1}{2} \varpi_* v^2 \right]$$

2) Single temperature equation

- Evolves sum over all species
- T_e/T_i constant throughout time, implicitly assuming
 - Instantaneous thermal equilibration
 - Split of losses between species evolves as pressure ratio changes

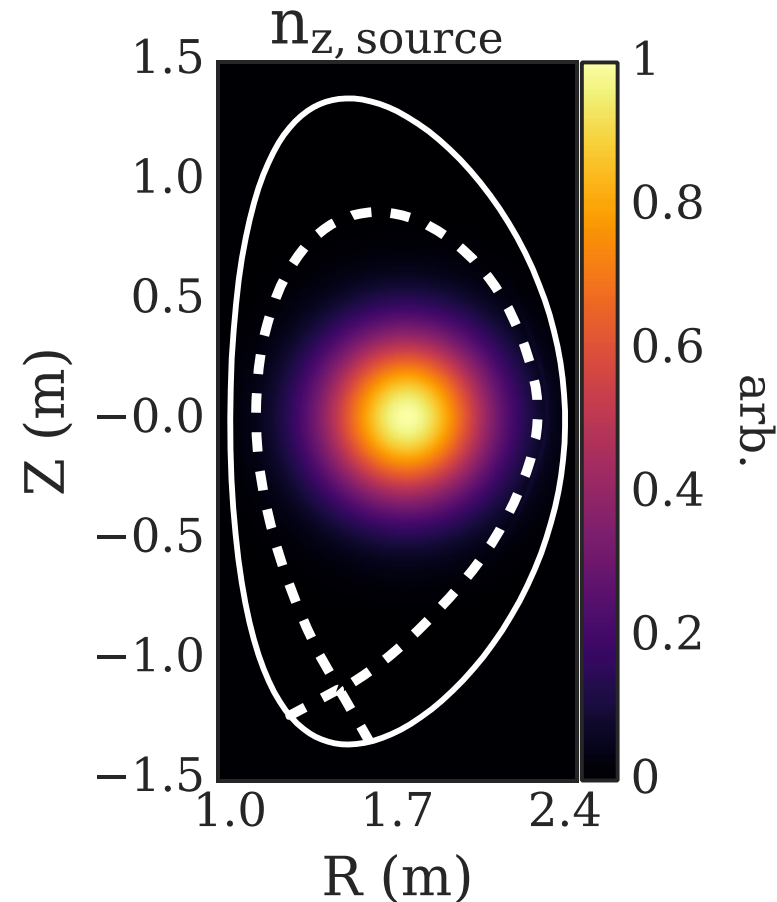
*N.M. Ferraro et al. Nucl. Fusion 59 016001 (2019).

3D Nonlinear Modeling of Core Impurity Injection

Fast Impurity Injection in DIII-D Core Used for Cross-Code Benchmarking with NIMROD

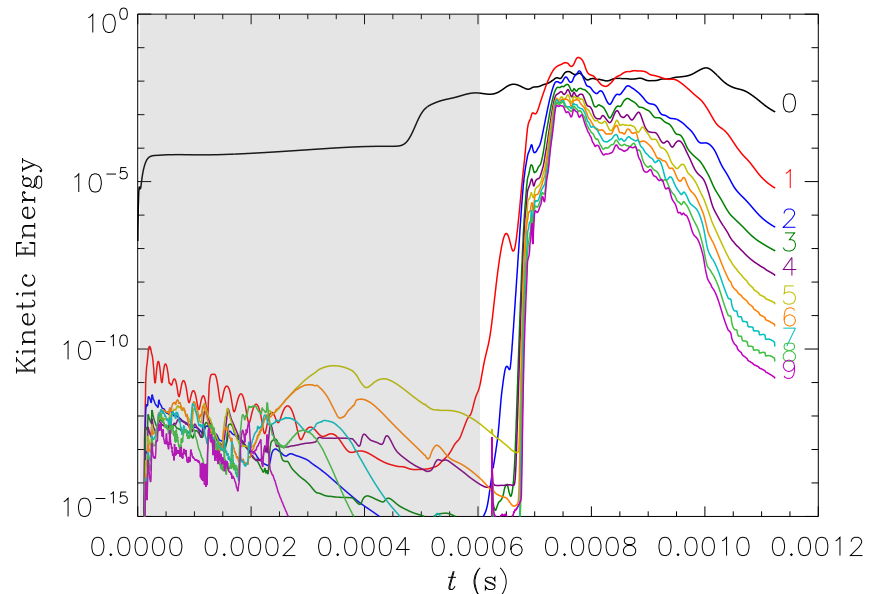
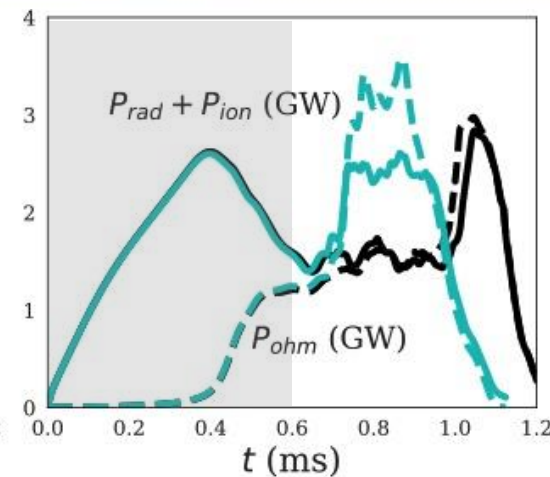
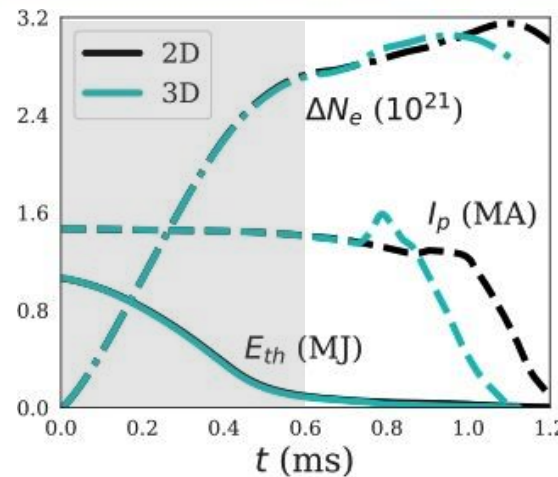
- **Simulation setup**
 - DIII-D shot 137611 @ 1950 ms
 - Single-fluid, single temperature
 - Fixed boundary
- **Continuous neutral impurity deposition (Ar or Ne)**
 - No impurities to start
 - Gaussian source
$$\frac{dn_z}{dt} = \nu \frac{R_0}{R} \exp \left[-\frac{(R - R_0)^2 + (Z - Z_0)^2}{2\delta^2} \right]$$
 - Injection rate ~1 mm Ne/Ar pellet per ms
- **Highly successful axisymmetric benchmark completed †**

† [B.C. Lyons et al., Plasma Phys. Control. Fusion, 61, 064001 \(2019\).](#)



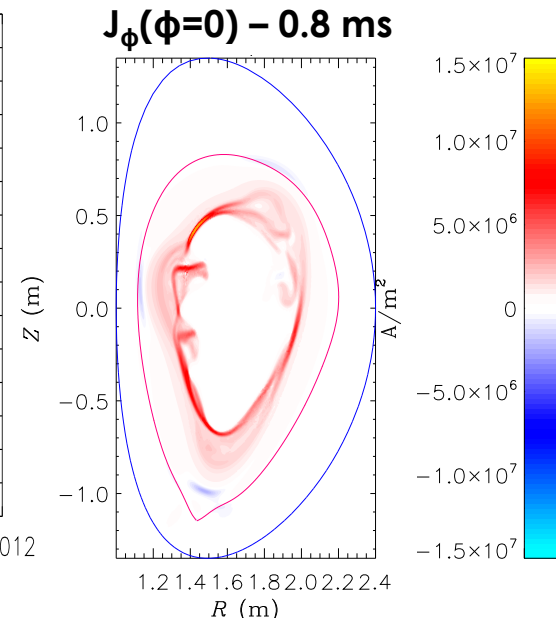
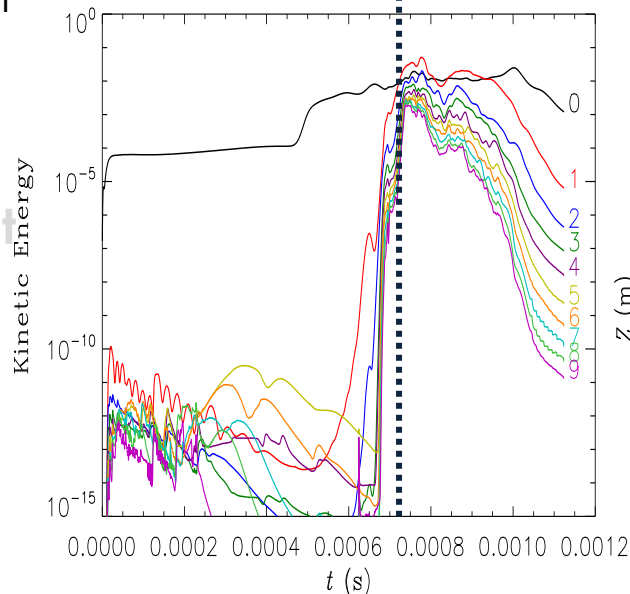
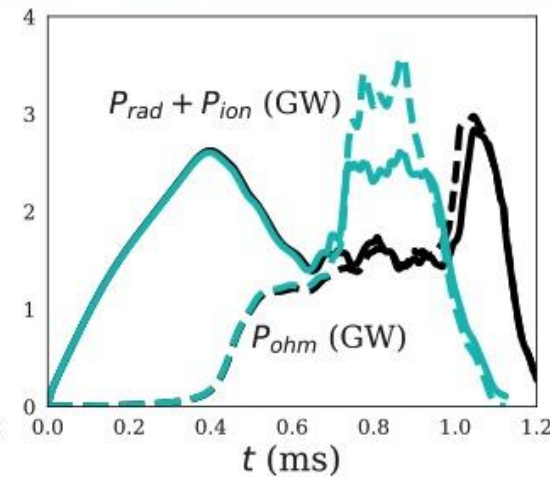
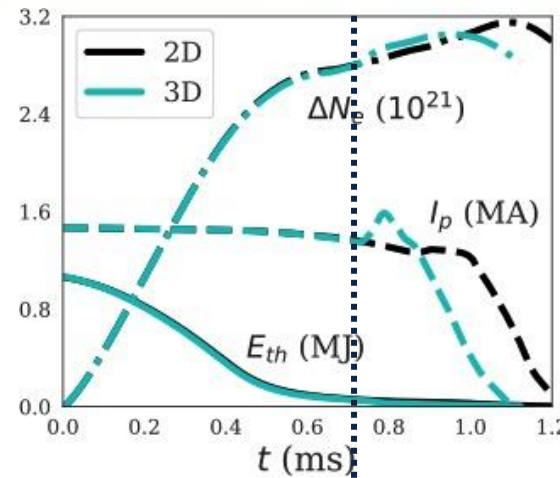
3D M3D-C1 Modeling Shows Stable Thermal Quench, Instability-Induced Current Quench with I_p Spike

- **3D, nonlinear simulation performed with M3D-C1 using argon benchmark initial conditions**
 - 3D run is linearly stable throughout thermal quench due to axisymmetric deposition
 - Plasma sheet goes unstable, quenching current
- **Instabilities cause current to spike**
 - Axisymmetric current broadens significantly
 - First spike of this magnitude in 3D MHD simulation



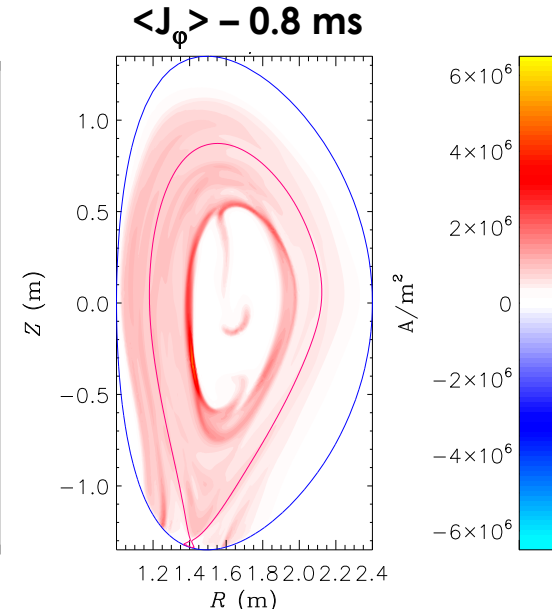
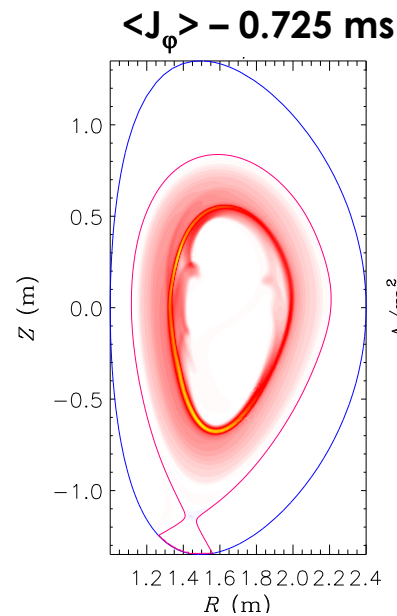
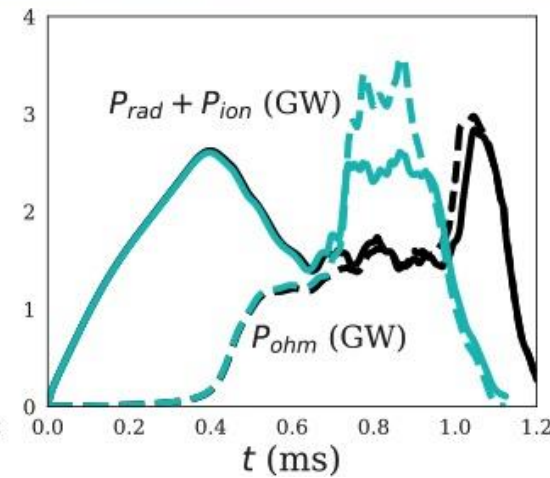
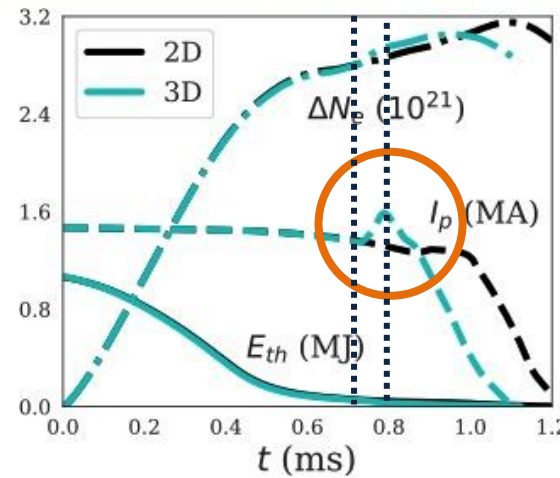
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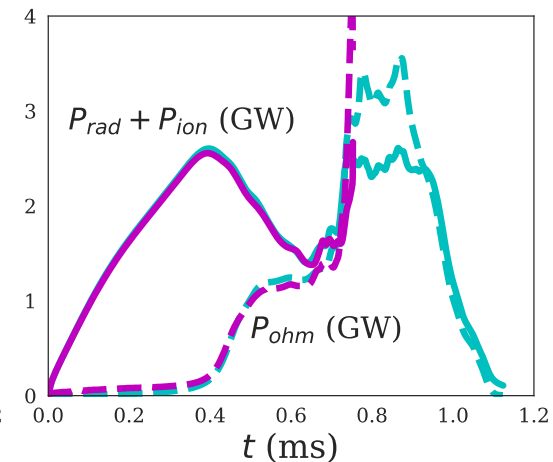
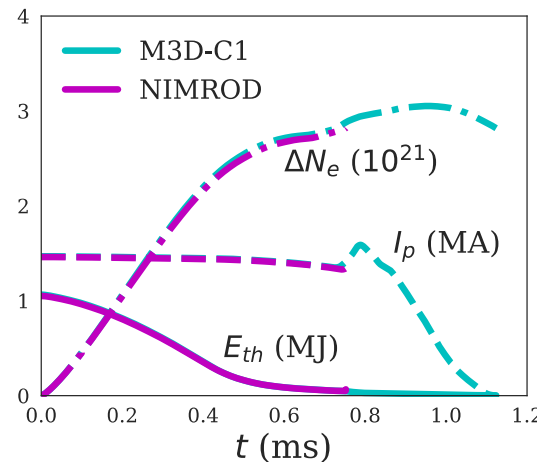
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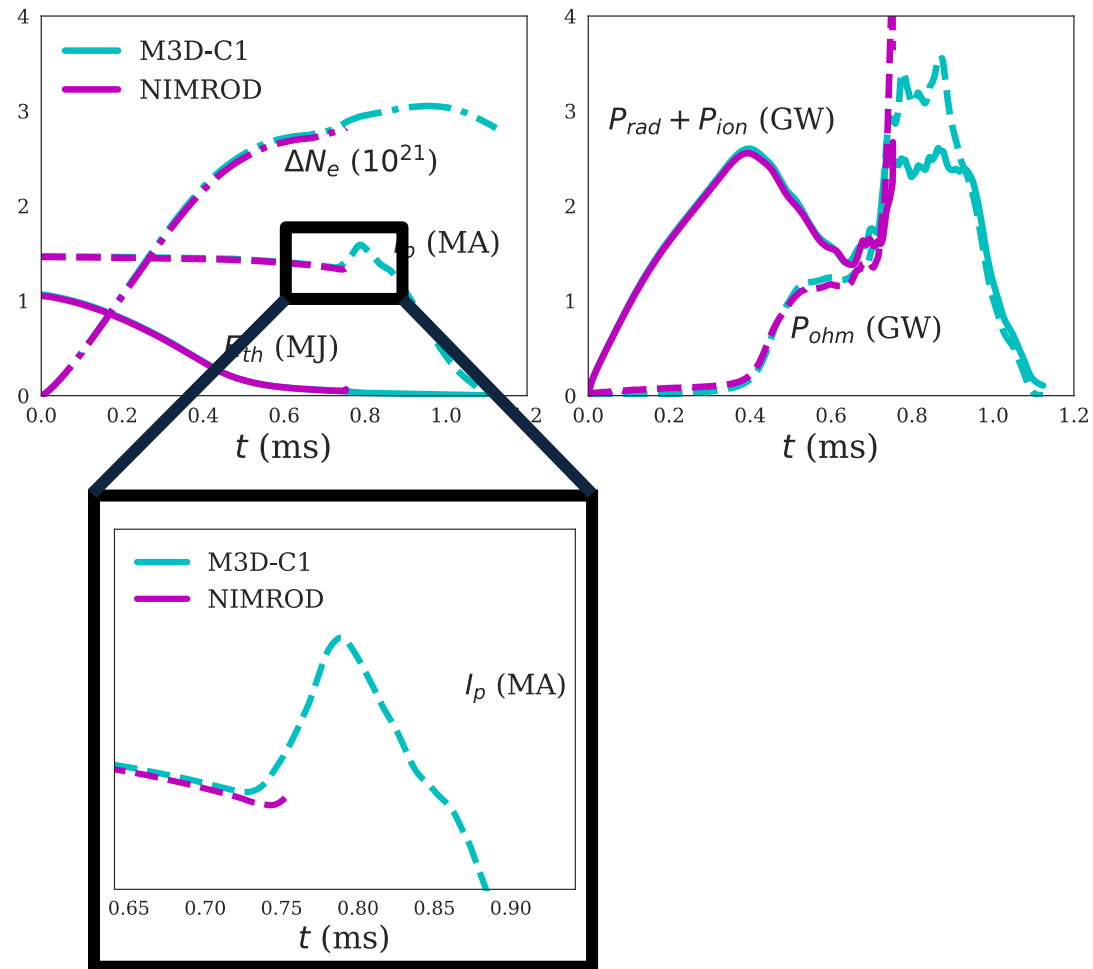
3D Benchmark with NIMROD Has Begun

- Shows axisymmetric behavior through ~0.75 ms, like M3D-C1
- Numerical instabilities hindering study of 3D crash
- Beginning to see current spike, though slightly delayed
- Benchmark with edge-injected ablating pellet also underway



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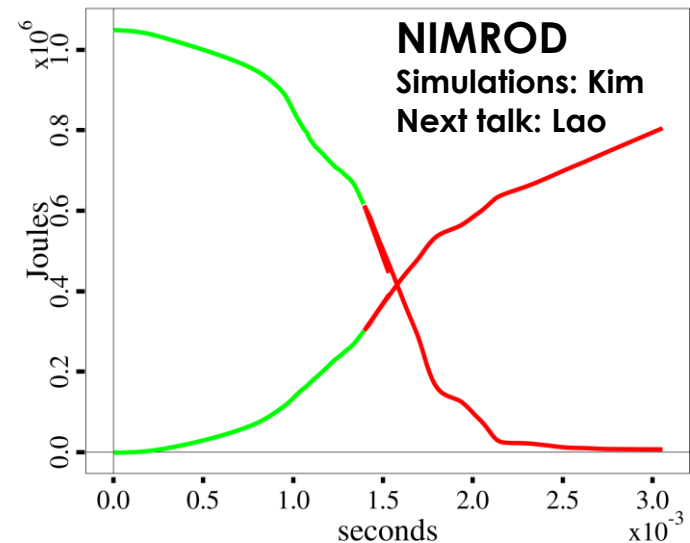
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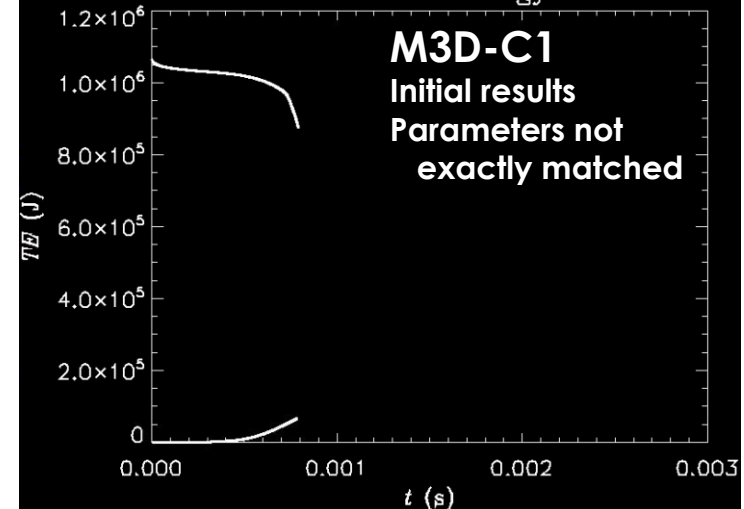
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Thermal and Radiated Energy vs. t



Thermal Energy



Pellet Mitigation Modeling

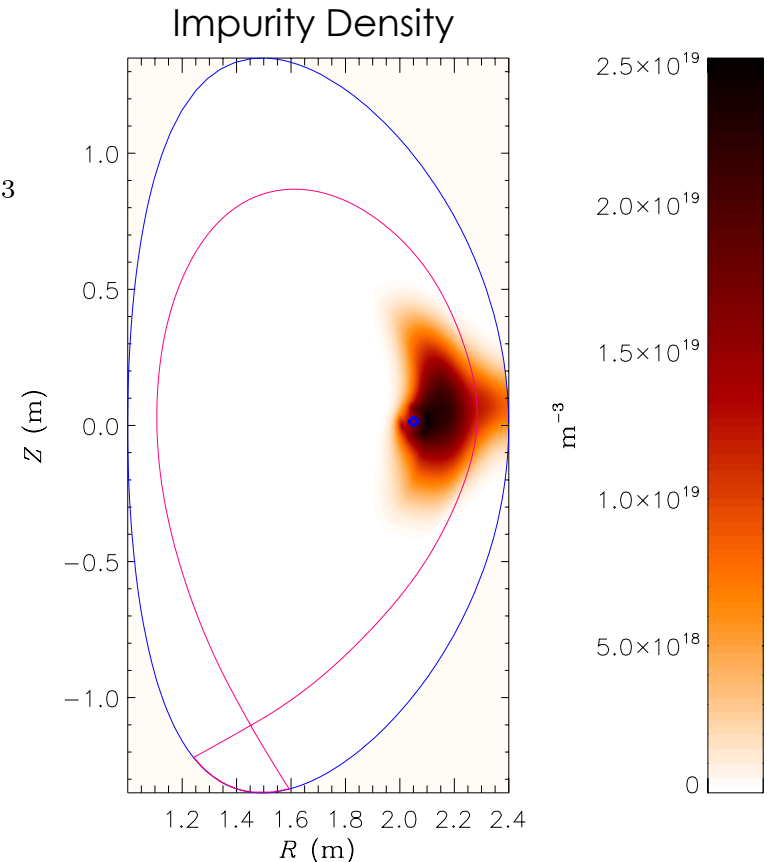
Ablation Model for Ne-D2 Pellets Implemented in M3D-C1

- **Practical, analytic expression fit to more complex ablation model (Parks)**

$$G(\text{g/s}) = \lambda(X) \left(\frac{T_e}{2000 \text{ eV}} \right)^{5/3} \left(\frac{r_p}{0.2 \text{ cm}} \right)^{4/3} \left(\frac{n_e}{10^{14} \text{ cm}^{-3}} \right)^{1/3}$$

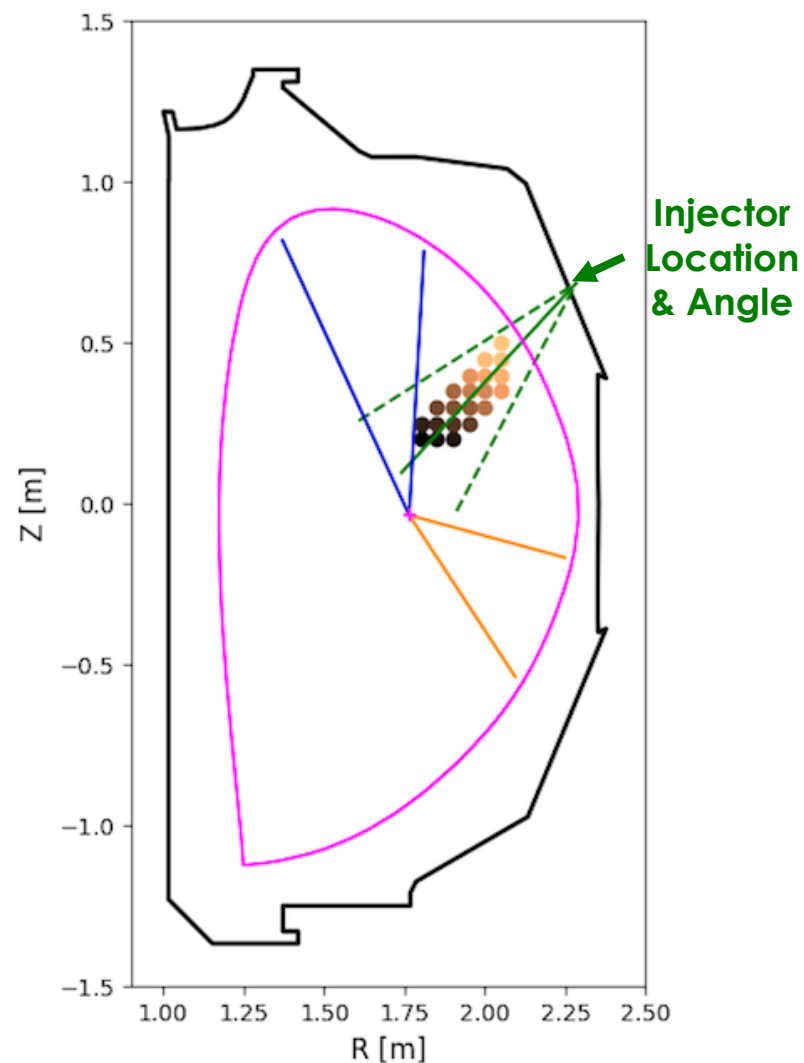
λ is fitting function, depending on molar X fraction of D2

- **M3D-C1 implementation**
 - Advance pellet location in time
 - Calculate number of particles ablated and pellet-surface recession at each time step
 - Deposit main ion and/or impurities onto arbitrary spatial distribution (e.g. 2D or 3D Gaussian)



Initial M3D-C1 DIII-D SPI Validation

- **Single, unshattered pellet modeled**
 - 4 mm radius
 - 200 m/s @ 41° from horizontal
 - Two compositions
 - Pure Ne
 - 1:10 Ne:D₂
- **Impurity deposition**
 - (R,Z) Gaussian with 5 or 10 cm half-width
 - von Mises distribution toroidally with 1.5 or 3 m half-width
- **Realistic density, temperature, and pressure profiles**
- **Single fluid velocity but separate T_e and T_i evolution**



Simulations Often Limited by Negative Temperatures and Numerical Instabilities

- **Negative temperatures can result due to competing desirable simulation parameters**
 - Low edge temperature to prevent heating of quenching plasma from the wall
 - Difficult to find good thermal conduction parameters
 - If $\kappa_{\parallel}/\kappa_{\perp}$ is large, steep edge gradients cause negative temperatures in SOL
 - If κ_{\parallel} is small, get cold spots around pellet that are difficult to resolve
 - If κ_{\perp} is large, significant thermal energy lost to diffusion
- **Temperature-evaluation floors improve stability, but runs still crash early**

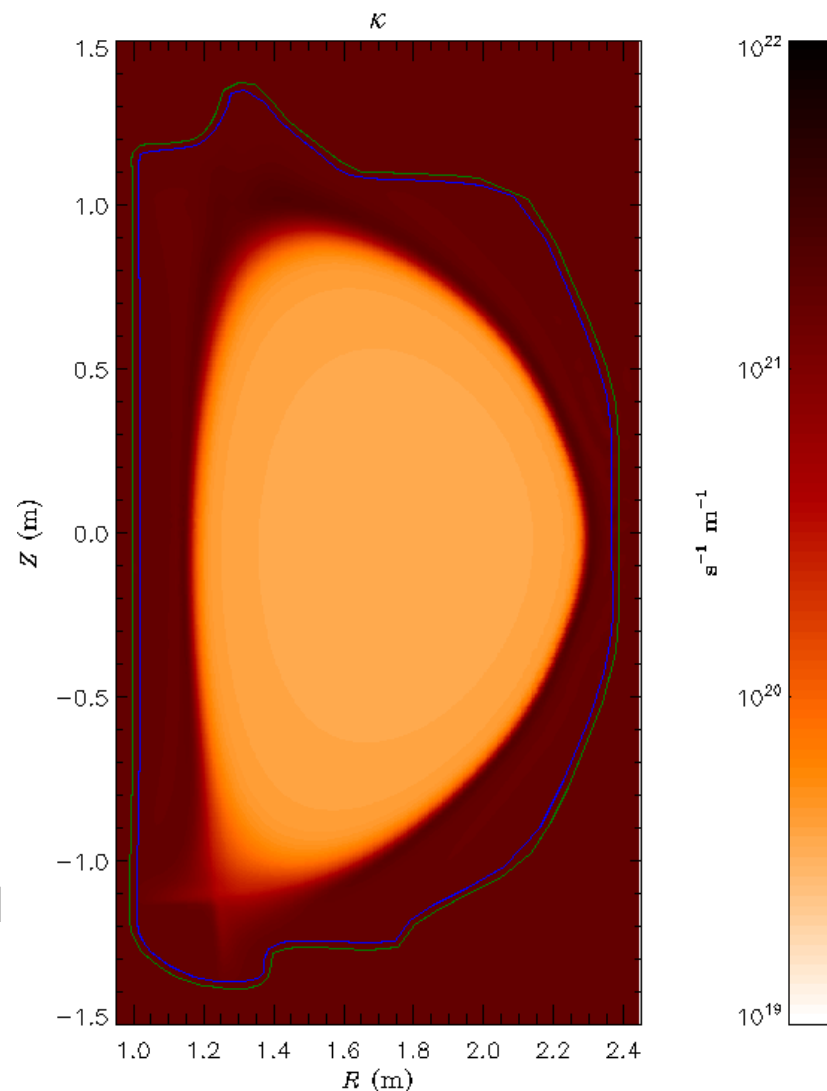
Temperature-Dependent Thermal Conductivity Improves Numerical Stability

- **Thermal conduction**

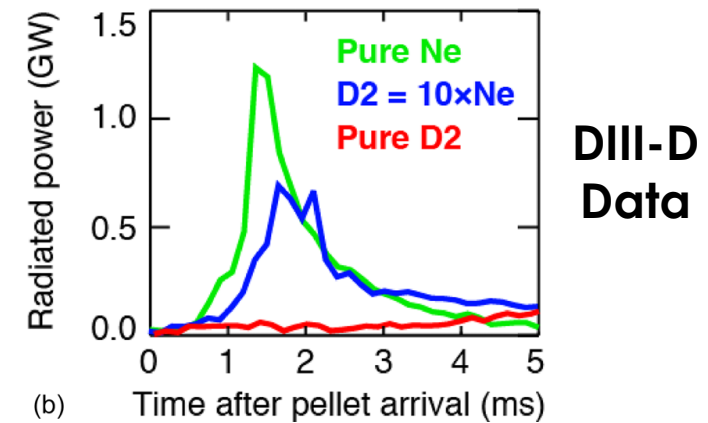
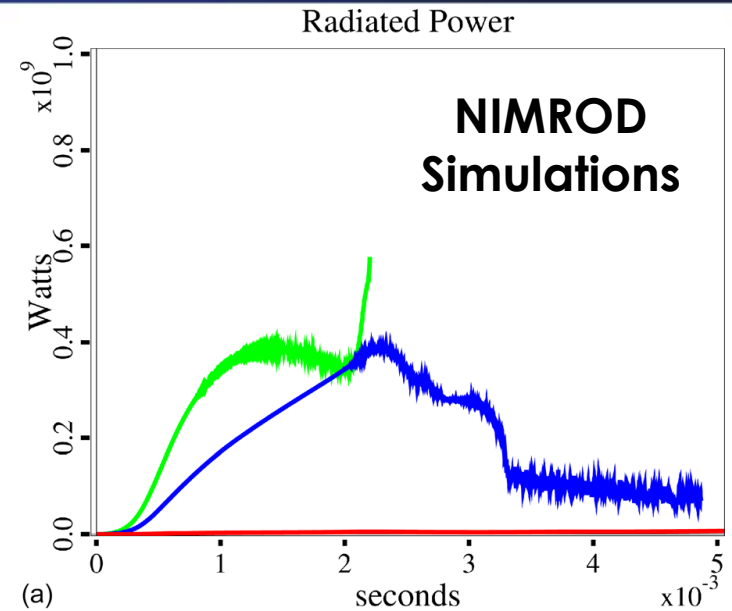
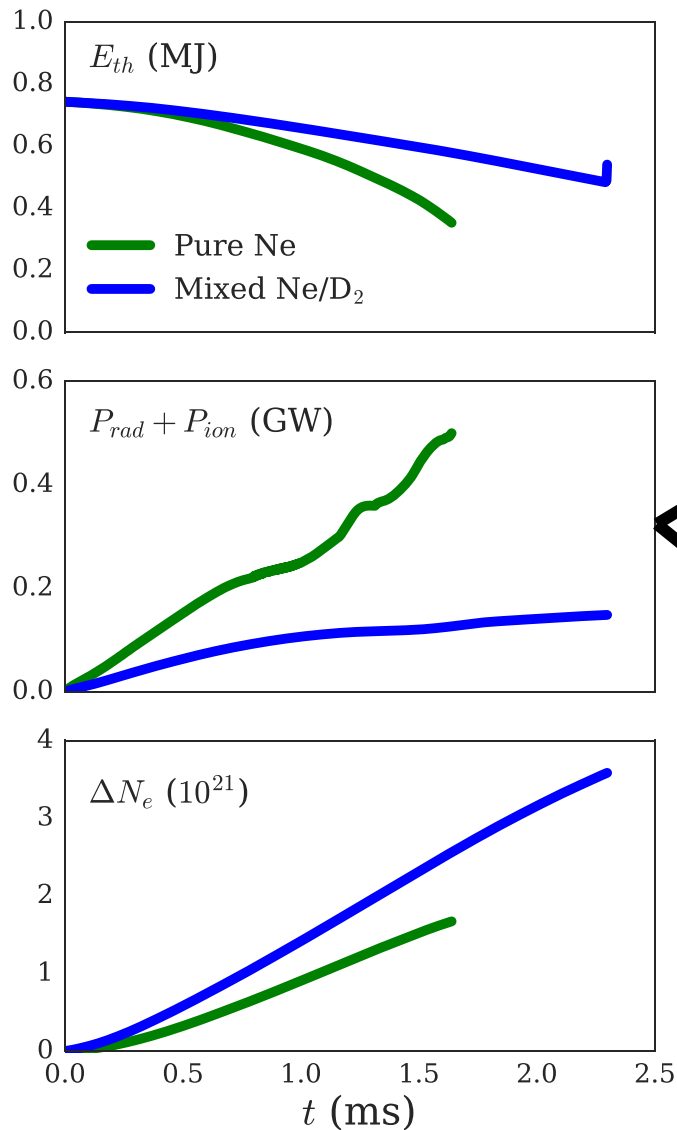
- Low, constant $\kappa_{\mathcal{R}}$ in core to maintain thermal energy during pellet injection $\kappa_{\mathcal{R}} = 3.33 \times 10^{19} \text{ m}^{-1}\text{s}^{-1}$
- Rises at low T_e as $1/T_e$ to maintain positivity during quench
- High, constant κ_{\parallel} to prevent strong gradients near pellet
 $\kappa_{\parallel}/\kappa_{\mathcal{R}} = 10^8$ in core

- **Other transport parameters**

- Realistic Spitzer resistivity ($S \sim 10^8$)
- Constant density diffusivity
- Constant viscosity, though rises in open-field-line region for numerical stability

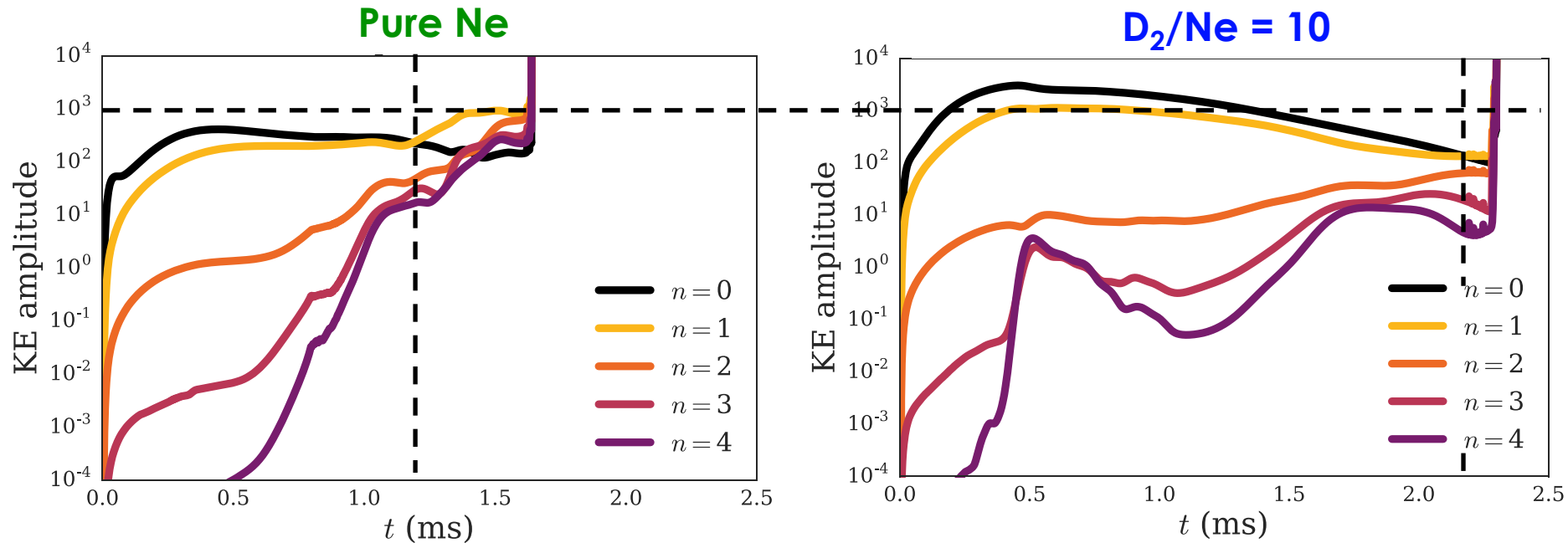


Pure Neon Pellet Radiates Thermal Energy Faster, Consistent with Experiment and NIMROD Modeling



C.C. Kim et al., Phys. Plasmas 26, 042510 (2019)

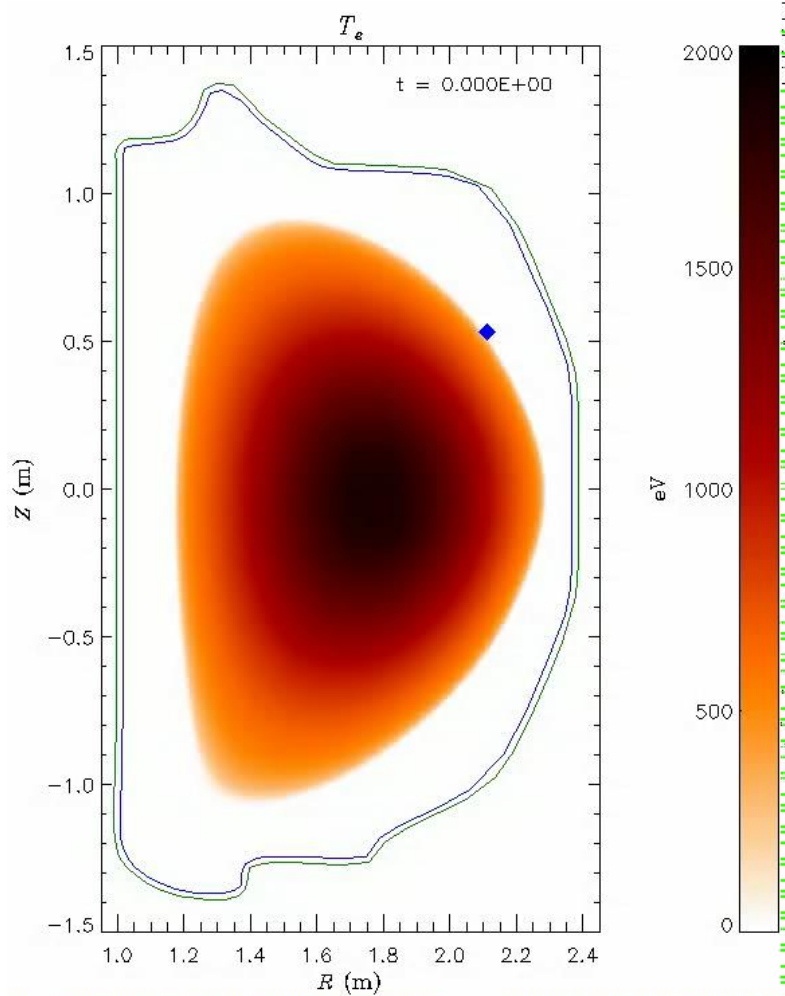
Pure Neon Pellet Induces Dominant $n=1$ Mode, Mixed Pellet Remains More Quiescent



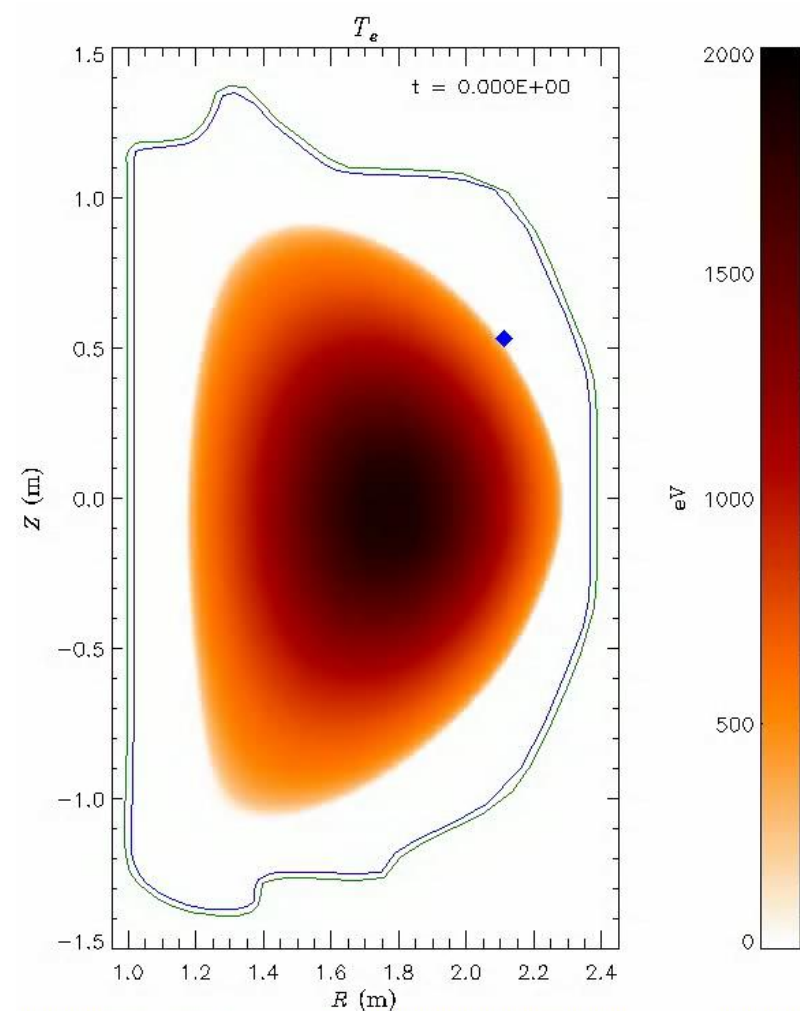
- **Late-time MHD appears marginally resolved toroidally**
- **Future work to extend simulations**
 - More planes
 - Smaller time step
 - More dissipation

Pure Neon Pellet Induces Dominant n=1 Mode, Mixed Pellet Remains More Quiescent

Pure Ne

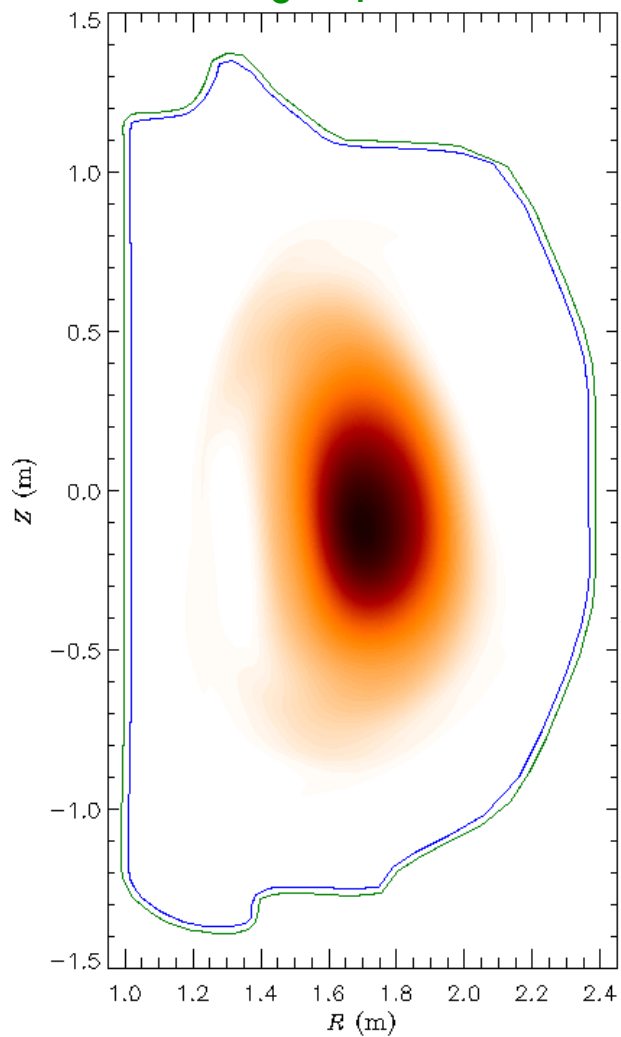


$D_2/Ne = 10$

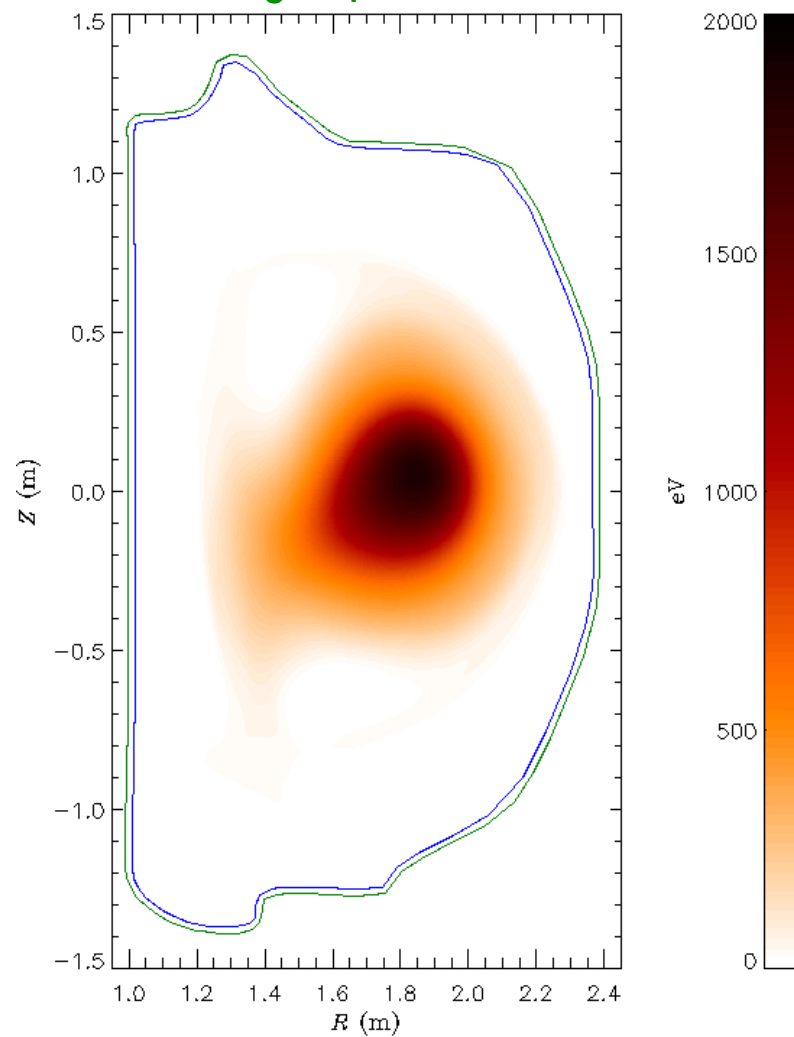


Pure Ne is Highly 3D at 1.5 ms

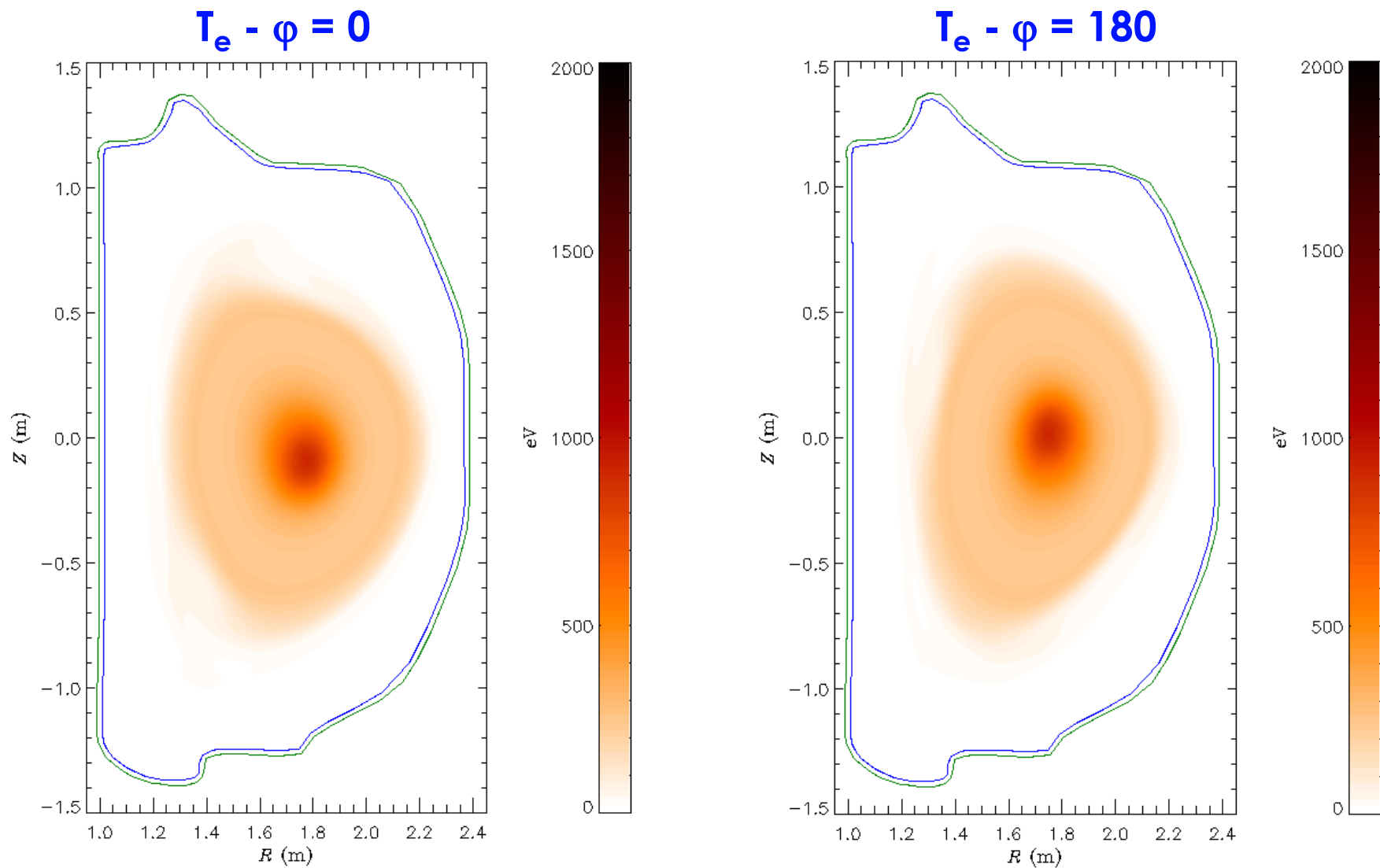
$T_e - \phi = 0$



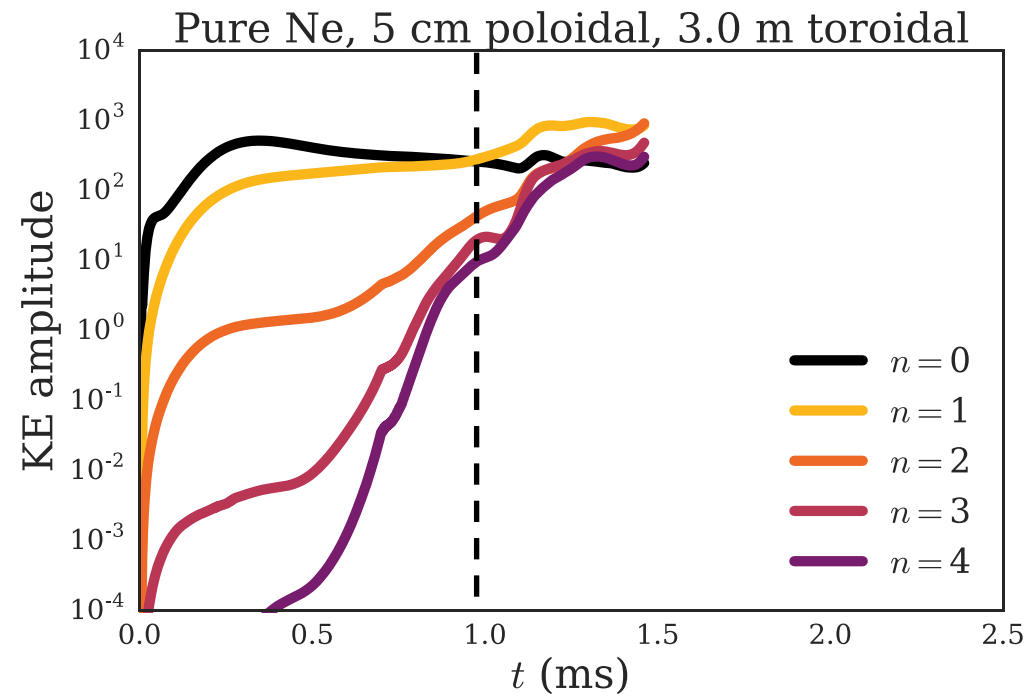
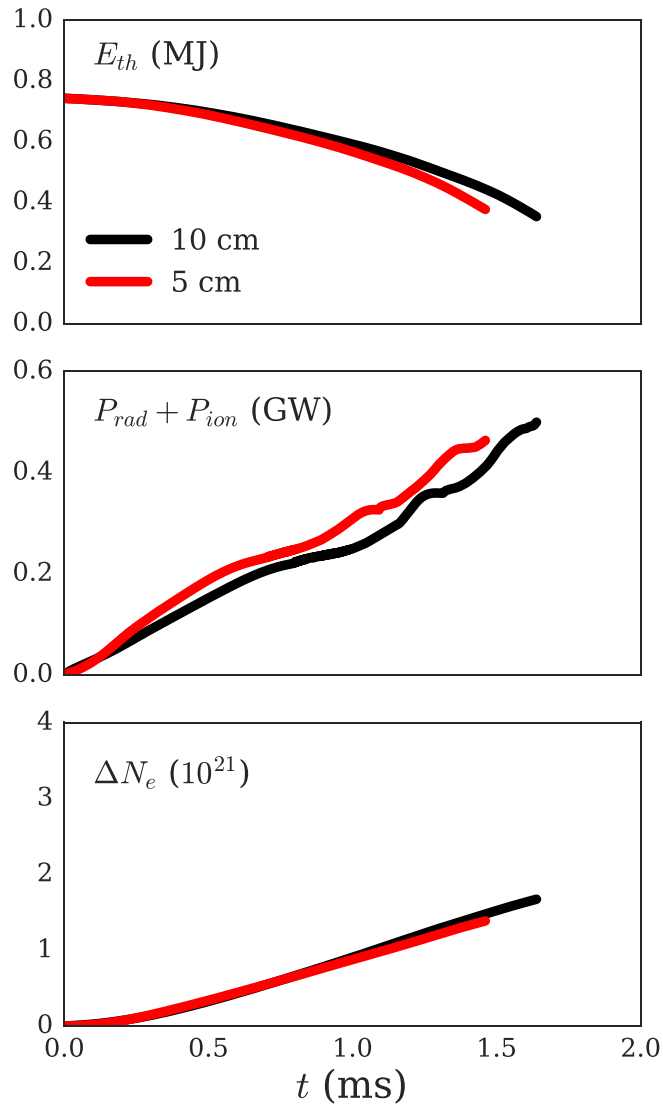
$T_e - \phi = 180$



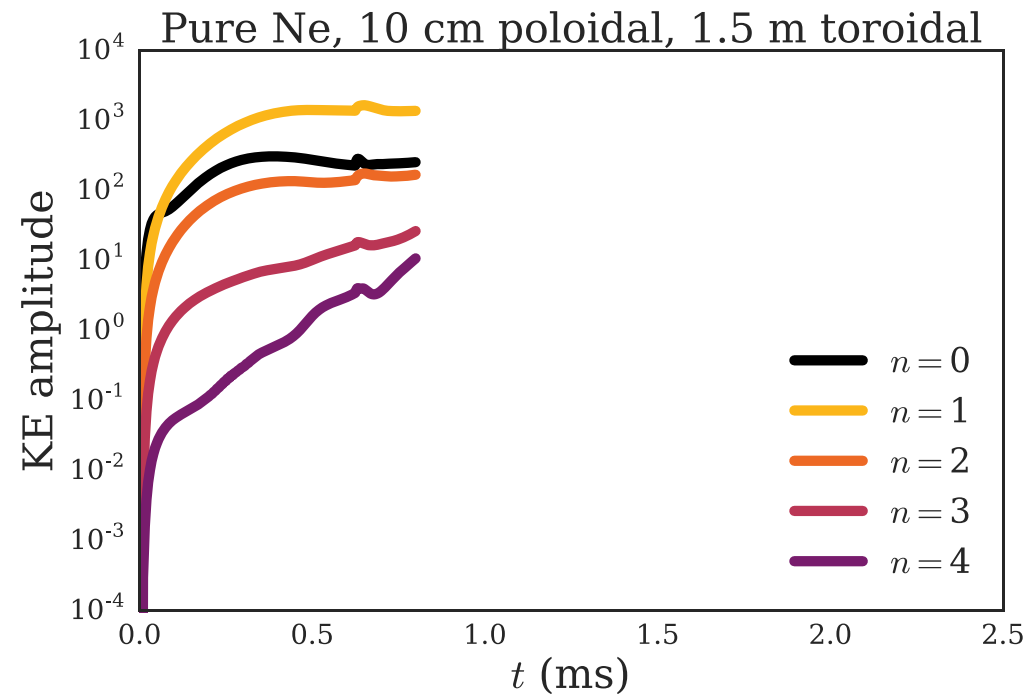
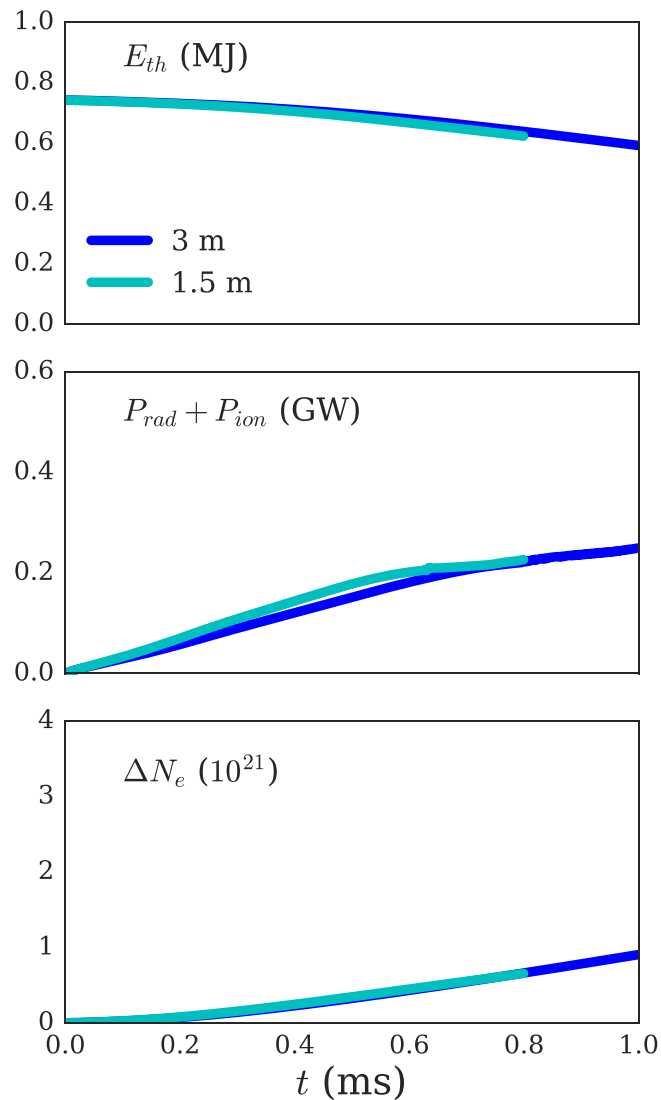
Mixed Ne/D₂ Remains More Axisymmetric through 2.25 ms



Poloidally Localized Pellet Drives Slightly Earlier MHD and More Radiation



Toroidal Localization Drives Early Non-Axisymmetry, but Not Instability or Radiation



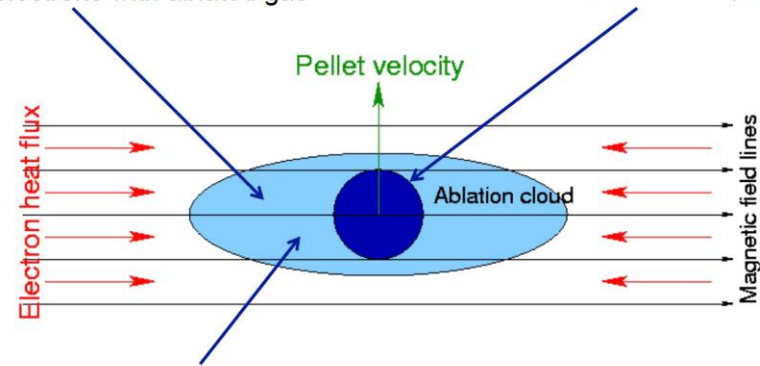
Coupling to Pellet Ablation Code

M3D-C1 is Coupling to Lagrangian-Particle Ablation Code is Underway

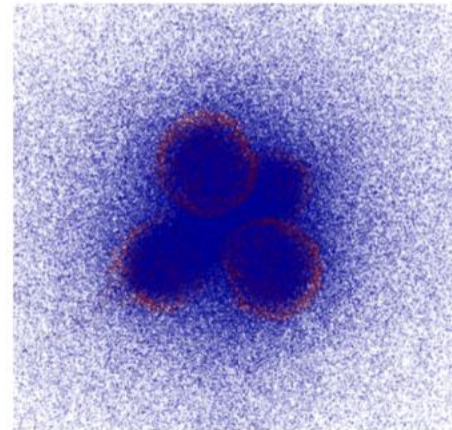
- **New Lagrangian particle code under development (Samulyak – Stonybrook)**
- **Resolves detailed physics**
 - MHD & atomic processes in ablation cloud
 - Phase transition at pellet/ablation surface
- **M3D-C1 and LP code exchange necessary information**
 - MHD code send upstream plasma parameters
 - Ablation code sends far-field ablated density

- Kinetic model for the interaction of hot electrons with ablated gas

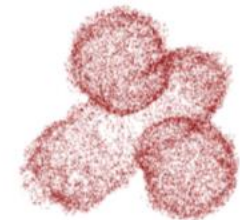
- Explicitly tracked pellet surface
- Phase transition (ablation model)



- Low Magnetic Re MHD equations
- Equation of state with atomic processes (Zeldovich average ionization model and tabular EOS based on solution of Saha equations)
- Radiation model
- Electric conductivity model



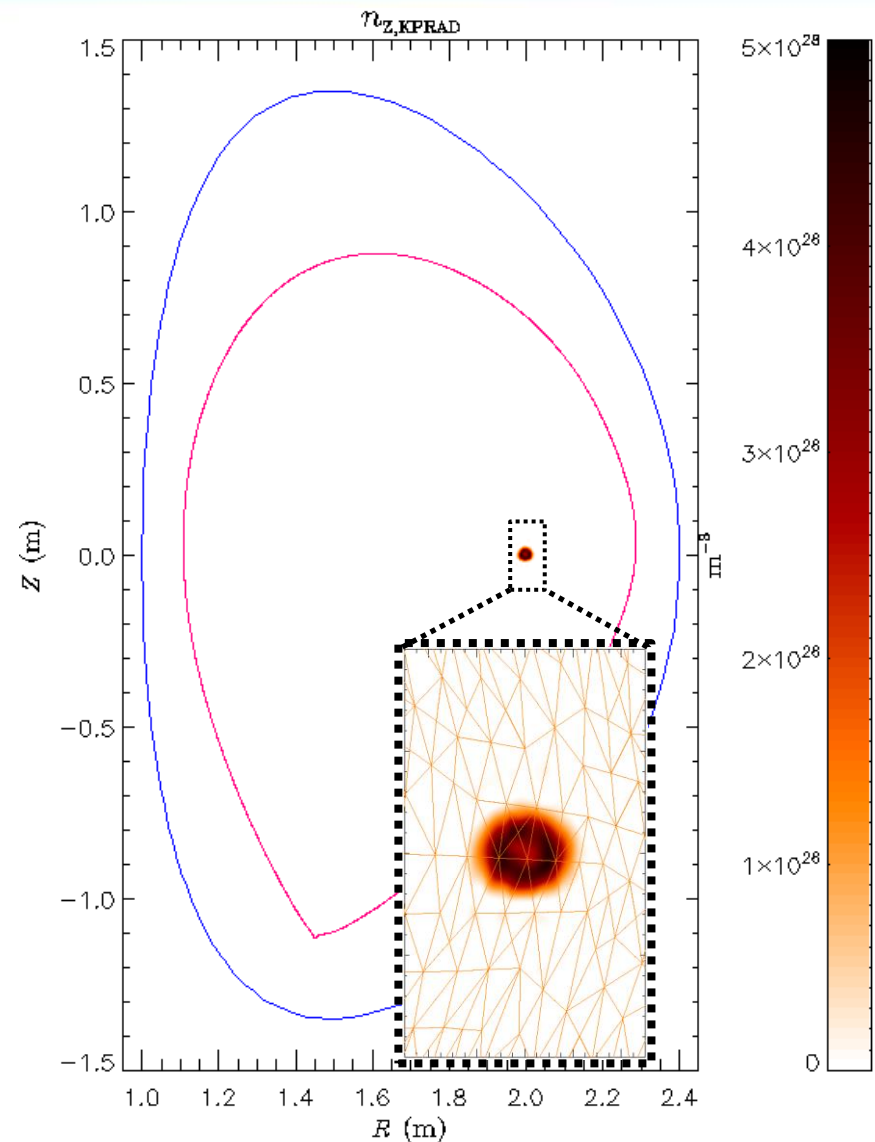
(a) SPI ablation cloud



(b) Particles to be sent to tokamak code

M3D-C1 Successfully Deposits Lagrangian Particles

- **LP code run for stationary 2 mm Ne pellet**
 - Magnetic field: 2 T
 - Electron density: 10^{20} m^{-3}
 - Electron temperature: 2 keV
- **Each LP of ablated material written to text file**
- **Read into M3D-C1**
 - Each particle is delta-function source
 - Interpolated onto finite-element mesh



Summary

- **M3D-C1 is now being used to simulate disruption mitigation by pellet injection in DIII-D**
- **3D modeling of benchmark case showed significant current spike driven by instabilities**
- **Initial SPI validation shows proper trends with pellet composition**
- **Increased localization of impurity deposition**
 - Poloidal: more MHD and radiation
 - Toroidal: more non-axisymmetry, but not instability
 - More challenging but likely more accurate
- **Ablated material from LP code successful deposited in M3D-C1**

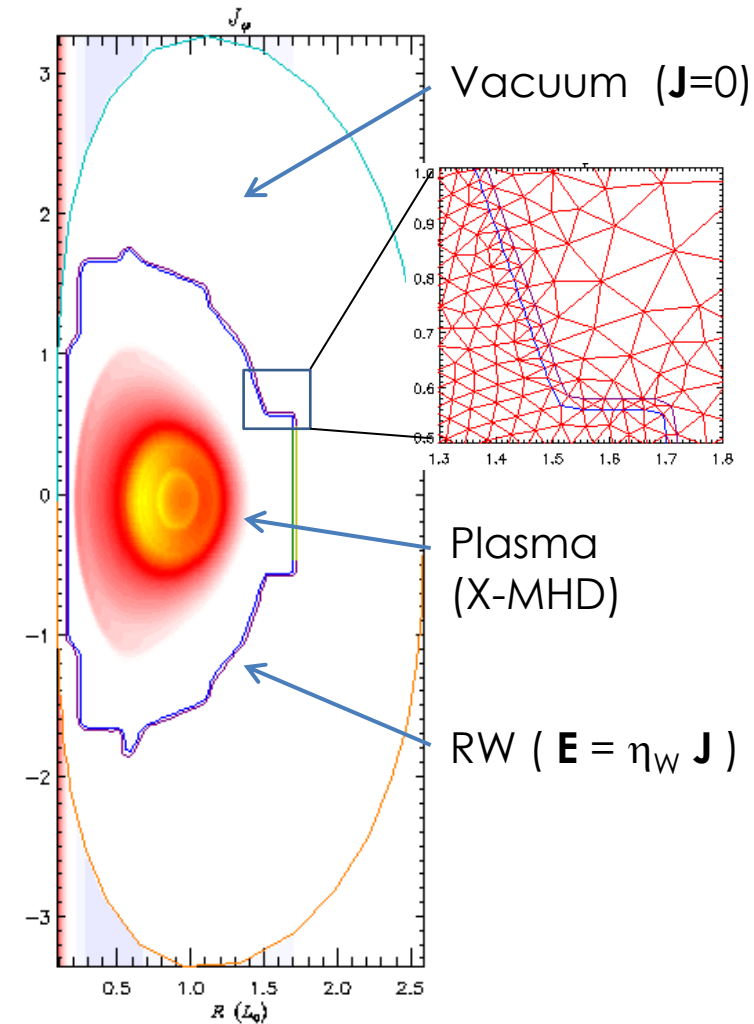
Future M3D-C1 Disruption Mitigation Work

- **Complete 3D nonlinear benchmark with NIMROD**
- **Continue validation against DIII-D experiments**
- **Continue study of toroidal localization of pellet source**
 - Axisymmetric vs. extended Gaussian vs. spherical
 - Make use of toroidal packing
- **Simulations with multiple pellet sources**
 - Shattered fragment cloud
 - Multiple toroidal injection (underway)
 - Couple to Lagrangian-particle/ablation code
- **Prediction and validation for JET, KSTAR, & ITER**
- **Dynamic simulations iterating between M3D-C1 and LP ablation code**

Additional Slides

M3D-C1* Solves the Extended-MHD Equations

- Three-dimensional toroidal geometry
- Full (not reduced) MHD
- Solves for potential and stream-function fields for \vec{A} & \vec{v} ($\nabla \cdot \vec{B} = 0$ intrinsically)
- Includes resistivity, density diffusivity, viscosity, & thermal conductivity
- Two-fluid effects (optional)
- 3D high-order finite elements
 - Unstructured, triangular mesh in poloidal plane
 - Structured toroidally, but can pack planes
- Can solve with finite-thickness resistive wall in domain**



*S. C. Jardin, et al., *Comput. Sci. Discovery* 5, 014002 (2012).

**N.M. Ferraro, et al., *Phys Plasma* 23 056114 (2016).

M3D-C1 Solves the Extended-MHD Equations

Blue terms are 2-fluid

$$\frac{\partial n}{\partial t} + \nabla \cdot (n\mathbf{V}) = \nabla \cdot D_n \nabla n + S_n$$

$$\frac{\partial \mathbf{A}}{\partial t} = -\mathbf{E} - \nabla \Phi, \quad \mathbf{B} = \nabla \times \mathbf{A}, \quad \mathbf{J} = \nabla \times \mathbf{B}, \quad \nabla_{\perp} \cdot \frac{1}{R^2} \nabla \Phi = -\nabla_{\perp} \cdot \frac{1}{R^2} \mathbf{E}$$

$$nM_i \left(\frac{\partial \mathbf{V}}{\partial t} + \mathbf{V} \cdot \nabla \mathbf{V} \right) + \nabla p = \mathbf{J} \times \mathbf{B} - \nabla \cdot \mathbf{\Pi}_i + \mathbf{S}_m$$

$$\mathbf{E} + \mathbf{V} \times \mathbf{B} = \frac{1}{ne} \left(\mathbf{R}_c + \mathbf{J} \times \mathbf{B} - \nabla p_e - \nabla \cdot \mathbf{\Pi}_e \right) - \frac{m_e}{e} \left(\frac{\partial \mathbf{V}_e}{\partial t} + \mathbf{V}_e \cdot \nabla \mathbf{V}_e \right) + \mathbf{S}_{CD}$$

$$\frac{3}{2} \left[\frac{\partial p_e}{\partial t} + \nabla \cdot (p_e \mathbf{V}) \right] = -p_e \nabla \cdot \mathbf{V} + \frac{\mathbf{J}}{ne} \cdot \left[\frac{3}{2} \nabla p_e - \frac{5}{2} \frac{p_e}{n} \nabla n + \mathbf{R}_c \right] + \nabla \cdot \left(\frac{\mathbf{J}}{ne} \right) : \mathbf{\Pi}_e - \nabla \cdot \mathbf{q}_e + Q_{\Delta} + S_{eE}$$

$$\frac{3}{2} \left[\frac{\partial p_i}{\partial t} + \nabla \cdot (p_i \mathbf{V}) \right] = -p_i \nabla \cdot \mathbf{V} - \mathbf{\Pi}_i : \nabla \mathbf{V} - \nabla \cdot \mathbf{q}_i - Q_{\Delta} + S_{iE}$$

$$\mathbf{V}_e = \mathbf{V}_i - \mathbf{J} / ne$$

$$\mathbf{R}_c = \eta ne \mathbf{J}, \quad \mathbf{\Pi}_i = -\mu \left[\nabla \mathbf{V} + \nabla \mathbf{V}^{\dagger} \right] - 2(\mu_c - \mu)(\nabla \cdot \mathbf{V}) \mathbf{I} + \mathbf{\Pi}_i^{GV}$$

$$\mathbf{q}_{e,i} = -\kappa_{e,i} \nabla T_{e,i} - \kappa_{\parallel} \nabla_{\parallel} T_{e,i}$$

$$\mathbf{\Pi}_e = (\mathbf{B} / B^2) \nabla \cdot \left[\lambda_h \nabla (\mathbf{J} \cdot \mathbf{B} / B^2) \right], \quad Q_{\Delta} = 3m_e (p_i - p_e) / (M_i \tau_e)$$

KPRAD Couples* to the M3D-C1 Pressure Equation(s)

3) Total and electron pressure equations

- Electrons lose energy to ionization and radiation
- Main ions lose energy only through cooling on electrons

$$\frac{\partial p}{\partial t} + \mathbf{v} \cdot \nabla p + \Gamma p \nabla \cdot \mathbf{v} = (\Gamma - 1) \left[\eta J^2 - \nabla \cdot \mathbf{q}_{tot} - \mathcal{P}_{rad} - \Pi_{tot} : \nabla \mathbf{v} + \frac{1}{2} \overline{\varpi} v^2 \right]$$

$$\frac{\partial p_e}{\partial t} + \mathbf{v} \cdot \nabla p_e + \Gamma p_e \nabla \cdot \mathbf{v} = (\Gamma - 1) \left[\eta J^2 - \nabla \cdot \mathbf{q}_e - \mathcal{P}_{rad} + Q_{ei} - \Pi_e : \nabla \mathbf{v} \right]$$

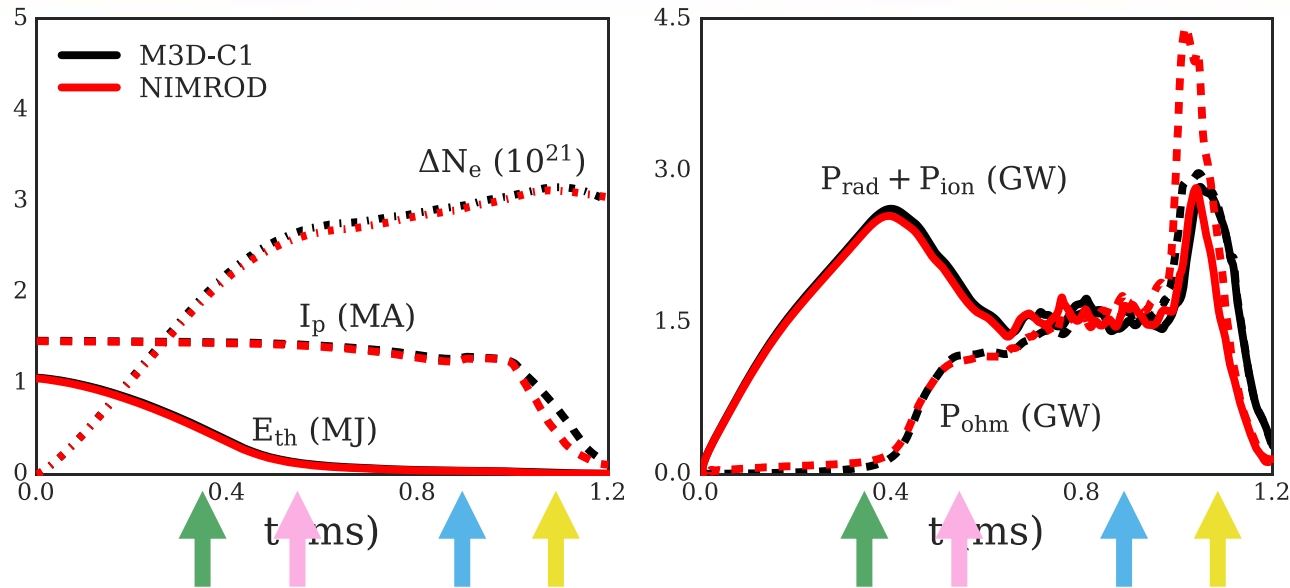
4) Single pressure equation

- Evolve only total pressure equation (above)
- p_e/p constant throughout time, implicitly assuming
 - No thermal equilibration
 - Losses split between ions and electrons by same fraction

*N.M. Ferraro et al. Nucl. Fusion 59 016001 (2019).

Successful Axisymmetric Benchmark between M3D-C1 and NIMROD*

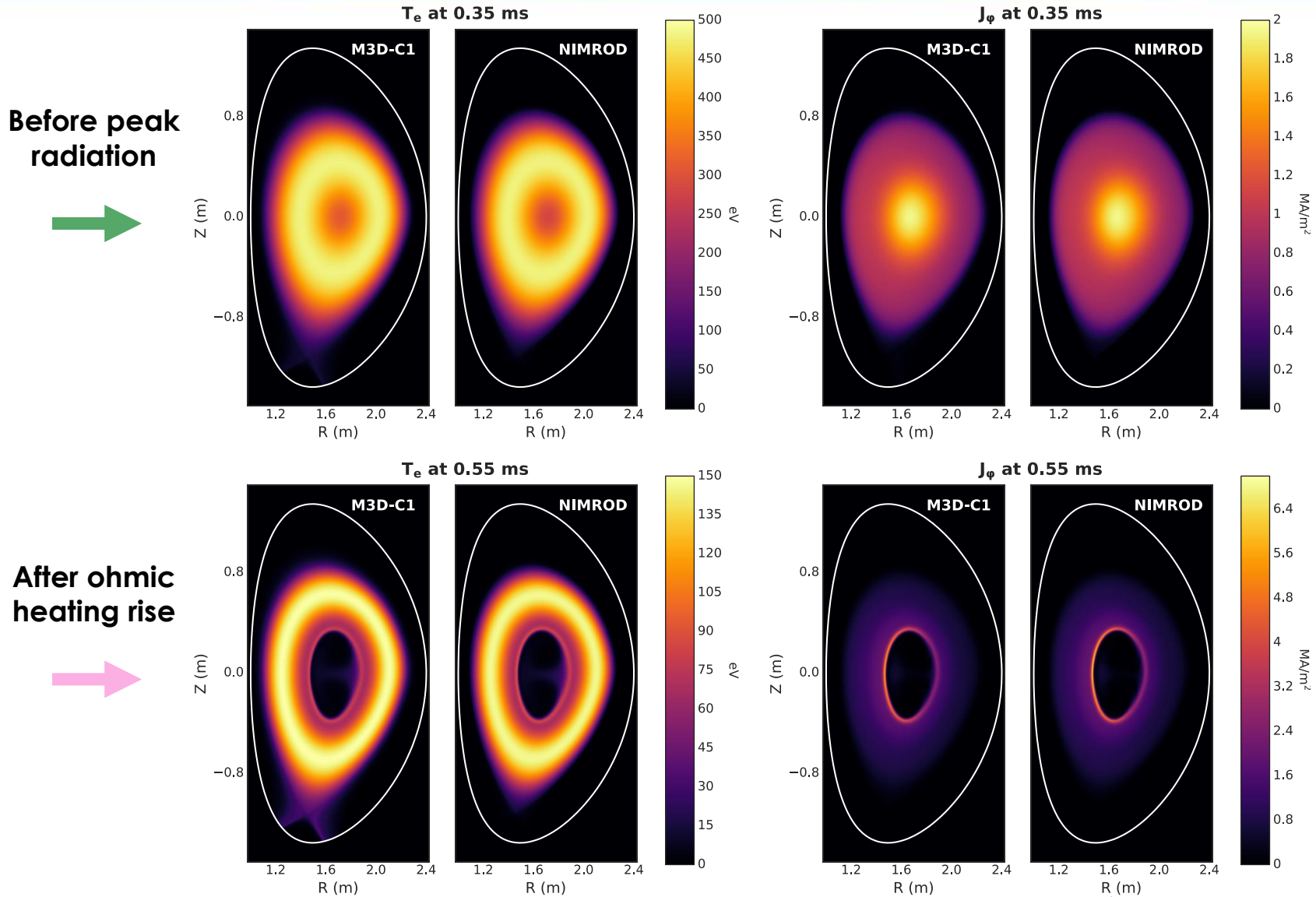
Argon Injection



- **Excellent agreement between codes**
 - Quantitative during thermal quench
 - Qualitative during current quench
- **Neon quench roughly 3x slower than argon**
- **Peak loss power when temp. on-axis falls near-zero**
- **Low temperature in core causes resistivity to rise**
 - P_{ohm} balances P_{loss}
 - Current drops more rapidly
- **Current quench caused by contact with boundary**
- **JOE^x benchmark underway**

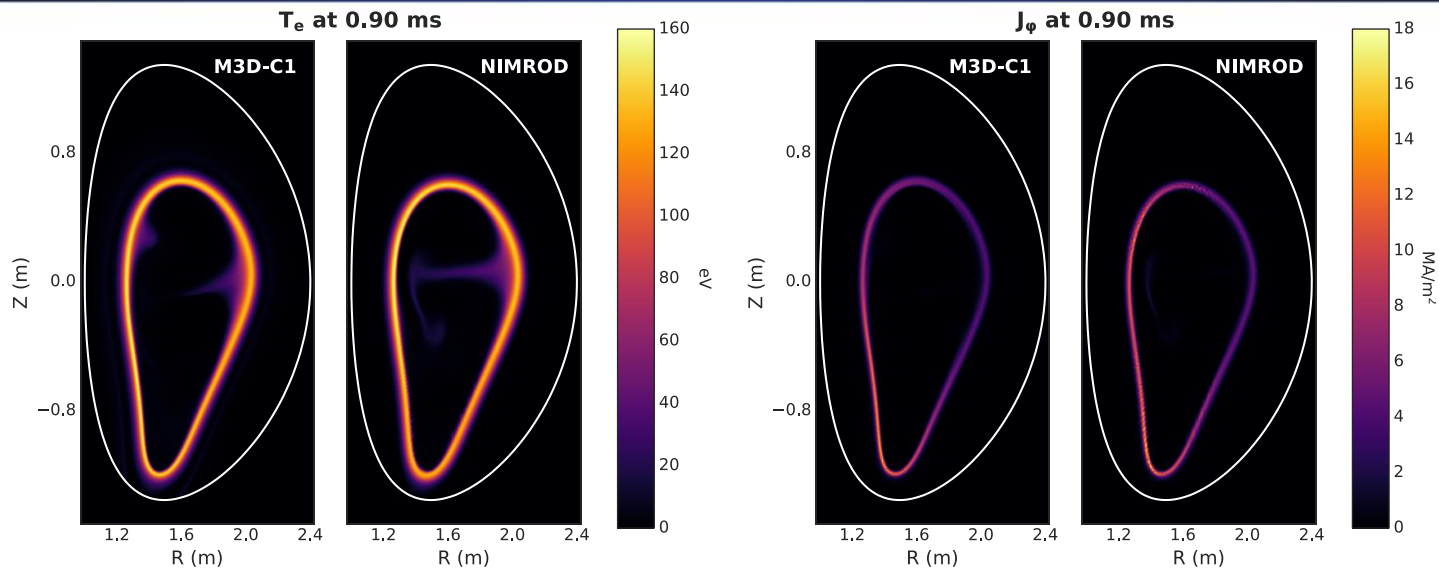
*C. R. Sovinec et al., J. Comput. Phys. 195, 355 (2004).
*G.T.A. Huysmans & O. Czarny. Nucl. Fusion 47, 659 (2007).

Early Times Show On-Axis Impurities Induce Inside-Out Thermal Quench and Hollowing of Current



Late Times Show Core Turbulence and Expanding Shell of Warm Plasma

$$P_{\text{ohm}} = P_{\text{loss}}$$



Rapid current
quench

