### Preconditioning over Multiple Fourier Components in NIMROD and NIMSTELL

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Center for Tokamak Transient Simulation Meeting October 16, 2022

Supported by US DOE grants DE-SC0018642 and DE-SC0018001

**Center for Tokamak Transient Simulation** 



### Motivation

- Nonlinear asymmetry leads to implicit coupling among toroidal Fourier harmonics in NIMROD.
- Toroidal shaping of the mesh and equilibria couple harmonics in NIMSTELL.
- Preconditioning over multiple Fourier components is important for efficiency.



Isosurfaces of  $J_{\parallel}/B$  from a nonlinear VDE computation with NIMROD.



Example |B| of 3D vacuum field that NIMSTELL solved on a twisting mesh.



# Many approaches to improve toroidal preconditioning have been tested in NIMROD.

- 1. Partial Gauss-Seidel-like updates with limited coupling within process layers helped in some cases.
- 2. Flexible GMRES (FGMRES) alternating between block Fourier diagonal (bFd) and block toroidal-plane diagonal is not better than bFd alone.
- 3. Kyle Bunkers tried Fourier multigrid and found that coarse-representation solves did not help the full solve.
- 4. 1D solves over toroidal angle showed little synergistic benefit with bFd with additive, multiplicative, or FGMRES approaches.
- 5. Fully coupled static condensation to reduce the system size reduces iteration count by ~ 50%.



# Direct solves over multiple Fourier components and the full poloidal plane has been implemented in NIMSTELL.

- Modifications to the *V*-matrix routine produces systems for bands of Fourier components.
  - Numerical integration is over all toroidal mesh points for FFTs.
  - Bands may overlap.
  - Systems use real algebra for separated real/imaginary components.





Example set of matrices with *stride* = 2 and *overlap* = 1.

## A parameter scan uses a cylindrical interchange mode in a twisted mesh.

- The m=4, n=1 mode is resonant at r = 0.371, where  $D_s = 0.443$ .
- Twisting the mesh one turn puts equilibrium  $B_{pol}$  into the n=1 component.



Profiles of safety factor (red) and Suydam parameter (black).

NIMSTELL linear result on  $A_z$ eigenfunction component with all dissipation parameters set to  $10^{-7}$ .



Twisting does not match the mode helicity.

# Thirteen computations with a small mesh compare performance when varying stride and overlap.

- Elements are bicubic. Fourier representation is  $0 \le n \le 5$ .
- Outer part of the mesh is a ring of 18x24 elements; 540 total elements.
- The twisted mesh leads to strong  $\pm 1$  coupling among Fourier components in the semiimplicit operator for  $\Delta V$ .
- Two computations use block Fourier diagonal ("bFd") preconditioning for the V-advance.
- Ten have varied stride ("st") and overlap ("ov") with real-algebra solves for the V-advance.
- The "st6 dg" uses our "diagonal" preconditioner applied to the full matrix.
- Computations were run on 5 cores of a shared-memory workstation without layer parallelism.
- Timings include the advances for **A**, *n*, and *T* that use bFd with SuperLU\_DIST (small contributions).



# Iteration is weakly dependent on stride; it jumps to low numbers with an overlap of 2.

- Computations are run 10 steps, and average iteration count for  $\Delta V$  is shown.
- Computations with overlap needed GMRES, which affects timings. Others used CG.





#### Comparing matrix and solve times shows a tradeoff.



• The reasonable performance of bFd-diagonal probably stems from the small problem size.



### Is a complex matrix for $n \ge 0$ possible?

- The working implementation uses real algebra with the system split into real and imaginary degrees of freedom.
- Our complex Fourier expansion uses:

$$A(\zeta) = A_0 + \sum_{n=1}^{N} \left[ A_n e^{in\zeta} + cc. \right] = A_0 + \sum_{n=1}^{N} \left[ 2A_{n_r} \cos(n\zeta) - 2A_{n_i} \sin(n\zeta) \right]$$

- A complex matrix for  $n \ge 0$  would use  $\frac{1}{2}$  of the memory.
- There are also fewer operations.
  - A test shows SLU\_DIST factoring time is reduced by 2.



#### A complex matrix implies a restriction.

- Real-algebra Fourier basis functions are  $\cos(n\zeta)$  and  $-\sin(n\zeta)$ .
- Complex-algebra Fourier basis functions are  $\exp(in\zeta)$ .
- If the system can be represented by a complex matrix,

$$\begin{pmatrix} a & b \\ c & d \end{pmatrix} \begin{pmatrix} x_1 \\ x_2 \end{pmatrix} = \begin{pmatrix} r_1 \\ r_2 \end{pmatrix} \rightarrow \begin{pmatrix} a_r + ia_i & b_r + ib_i \\ c_r + ic_i & d_r + id_i \end{pmatrix} \begin{pmatrix} x_{1_r} + ix_{1_i} \\ x_{2_r} + ix_{2_i} \end{pmatrix} = \begin{pmatrix} r_{1_r} + ir_{1_i} \\ r_{2_r} + ir_{2_i} \end{pmatrix}$$

then the equivalent real matrix is restricted:

$$\begin{pmatrix} a_{r} & -a_{i} & b_{r} & -b_{i} \\ a_{i} & a_{r} & b_{i} & b_{r} \\ c_{r} & -c_{i} & d_{r} & -d_{i} \\ c_{i} & c_{r} & d_{i} & d_{r} \end{pmatrix} \begin{pmatrix} x_{1_{r}} \\ x_{1_{i}} \\ x_{2_{r}} \\ x_{2_{i}} \end{pmatrix} = \begin{pmatrix} r_{1_{r}} \\ r_{1_{i}} \\ r_{2_{r}} \\ r_{2_{i}} \end{pmatrix}$$



#### A complex implementation indicates limitations.

- The complex operation is effective for the **V**-matrix equation with non-twisting meshes.
  - The Krylov method takes only 1 or 2 steps when including all Fourier components in the matrix.
- It does not converge with a twisting mesh.
  - Printing some components of the *real-algebra* version shows that real and imaginary parts of diagonals are not the same.
  - This violates the restriction on the previous page!



#### A simple example shows the limitation.

• Consider diagonal matrix elements that result from a diffusion operation

with diffusivity 
$$D(\zeta) = \begin{cases} D_0, \frac{3\pi}{4} \le \zeta \le \frac{5\pi}{4} \\ 0, \text{ otherwise} \end{cases}$$
  
• "Real" phase,  $\cos \zeta : D_0 \int_{\frac{3\pi}{4}}^{\frac{5\pi}{4}} d\zeta \left(\frac{d}{d\zeta}\cos\zeta\right)^2 = D_0 \left(\frac{\pi}{4} - \frac{1}{2}\right)$   
• "Imaginary" phase,  $-\sin \zeta : D_0 \int_{\frac{3\pi}{4}}^{\frac{5\pi}{4}} d\zeta \left(\frac{d}{d\zeta}\sin\zeta\right)^2 = D_0 \left(\frac{\pi}{4} + \frac{1}{2}\right)$   
• Complex representation,  $\exp i\zeta : D_0 \int_{\frac{3\pi}{4}}^{\frac{5\pi}{4}} d\zeta \left(\frac{d}{d\zeta}e^{i\zeta}\right) \left(\frac{d}{d\zeta}e^{-i\zeta}\right) = D_0 \frac{\pi}{2}$ 

• This example indicates that we must separate the phases in real matrices.

#### MPI communication for decomposition over $\boldsymbol{\zeta}$ is now complete.

- Matrix construction subdivides NIMROD's grid blocks.
  - Integration over all  $\zeta$  is needed.
  - The decomposition is the same as used for projecting vectors.
  - Asynchronous communication is used to gather contributions.
- Preconditioning operations need to share Fourier components.
  - Matrix-band harmonics may differ from layer decomposition.
  - Asynchronous communication is used to exchange harmonics.



#### Static condensation over 3D can be used in other ways.

• Static condensation is element-based matrix partitioning,

$$\left(\underline{M}_{11} - \underline{M}_{12}\underline{M}_{22}^{-1}\underline{M}_{21}\right)\underline{x}_1 = \underline{b}_1 - \underline{M}_{12}\underline{M}_{22}^{-1}\underline{b}_2 ,$$

where  $\underline{x}_2 = \underline{M}_{22}^{-1}(\underline{b}_2 - \underline{M}_{21}\underline{x}_1)$  has the degrees of freedom from element interiors.

- In the banded systems discussed above, the static condensation is over each band.
- Forming  $\underline{M}_{12}$ ,  $\underline{M}_{21}$ , and  $\underline{M}_{22}$ , involves element-local FE computations, which is straightforward for layer parallelism.
- Will the block-Fourier diagonals of the Schur complement (top equation) make a scalable preconditioner in combination with the 3D static condensation?



### **Conclusions and Discussion**

- Direct solves over multiple Fourier components have been implemented.
  - Full matrix computation and factoring will not be tractable for large problems.
- Overlapping 2 Fourier components shows good performance in the test case.
  - This may be related to the strong  $\pm 1$  coupling.
  - CG with additive overlap should be possible.
- Parallel domain decomposition has just been implemented.
- 3D static condensation with bFd is worth trying.
- The developments will be ported to NIMROD.

