

Continuum closure plans and efforts*

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CTTS work at USU

Continuum energetic particle physics implemented using δf drift kinetic equation (DKE) with stress tensor coupling in NIMROD's flow evolution equation.

▶ Resistive Wall Modes

- ▶ Piggyback on Tech-X resistive wall development in NIMROD.
- ▶ Yr 4 Explore kinetic effects on stability with NIMROD, M3D-C1/DK4D: Compare with MARS-K
- ▶ Yr 5 Explore disruptions caused by energetic particles (fishbone modes) interacting with RWMs.

Continuum Chapman-Enskog-like (CEL)-DKEs implemented with tight coupling of collisional friction, anisotropic stress, conductive heat flow and collisional energy exchange closures in NIMROD's fluid model.

▶ Neoclassical Tearing Modes

- ▶ Ys 1-2 Identify suitable equilibria for NTM/locked mode disruptions on DIII-D for modeling (Tech-X).
- ▶ Yr -1 Implement Ramos-form of CEL-DKE closures in NIMROD (USU).
- ▶ Yr 1 NIMROD DKE closures in fixed magnetic island geometry (USU) by testing anisotropic thermal transport with tight CEL-DKE/temperature coupling.

Various forms for the δf DKE

The drift-kinetic equation may be written using several sets of variables:

1. Hazeltine*

$$\bar{f}(U, \mu, \mathbf{x}, t) \rightarrow \frac{\partial \bar{f}}{\partial t} + \mathbf{v}_{gc} \cdot \nabla \bar{f} + \frac{dU}{dt} \frac{\partial \bar{f}}{\partial U} + \frac{d\mu}{dt} \frac{\partial \bar{f}}{\partial \mu} = C$$

2. Ramos†

$$\bar{f}(v_{\parallel}, v_{\perp}, \mathbf{x}, t) \rightarrow \frac{\partial \bar{f}}{\partial t} + \mathbf{v}_{gc} \cdot \nabla \bar{f} + \frac{dv_{\parallel}}{dt} \frac{\partial \bar{f}}{\partial v_{\parallel}} + \frac{dv_{\perp}}{dt} \frac{\partial \bar{f}}{\partial v_{\perp}} = C$$

3. NIMROD‡

$$\bar{f}(s, \xi, \mathbf{x}, t) \rightarrow \frac{\partial \bar{f}}{\partial t} + \mathbf{v}_{gc} \cdot \nabla \bar{f} + \frac{ds}{dt} \frac{\partial \bar{f}}{\partial s} + \frac{d\xi}{dt} \frac{\partial \bar{f}}{\partial \xi} = C$$

*Hazeltine and Meiss, Plasma Confinement (Adisson-Wesley, RedwoodCity, 1992)

†Ramos, Phys Plasmas 15, 082106 (2008)

‡Held, *et al*, Phys Plasmas 22, 032511 (2015)

δf DKE in NIMROD for studies of energetic particle effects

The δf DKE implemented in NIMROD uses the coordinates (\mathbf{x}, t, s, ξ) , where $s = v/v_0$ and $\xi = v_{\parallel}/v$

$$\begin{aligned}
 \mathbf{v}_{\text{gc}} &= v_0 s \xi \mathbf{b} + \frac{\mathbf{E} \times \mathbf{B}}{B^2} + \frac{T_0 s^2}{qB} (1 + \xi^2) \mathbf{b} \times \nabla \ln B \\
 &\quad + \frac{2T_0 s^2}{qB^2} \left[\xi^2 (\mathbf{I} - \mathbf{b}\mathbf{b}) + \frac{1}{2} (1 - \xi^2) \mathbf{b}\mathbf{b} \right] \cdot \nabla \times \mathbf{B} + \frac{mv_0 s \xi}{qB^2} \mathbf{b} \times \frac{\partial \mathbf{B}}{\partial t} \\
 \dot{s} &= -s \frac{d \ln v_0}{dt} + \frac{s(1 - \xi^2)}{2} \frac{\partial \ln B}{\partial t} + \frac{q}{2T_0 s} (\mathbf{v}_{\parallel} + \mathbf{v}_c) \cdot \mathbf{E} + \frac{s}{2} (1 + \xi^2) \frac{\mathbf{E} \times \mathbf{B}}{B^2} \cdot \nabla \ln B \\
 \dot{\xi} &= \frac{1 - \xi^2}{2\xi} \left\{ -\xi^2 \frac{\partial \ln B}{\partial t} + (\mathbf{v}_{\parallel} + \mathbf{v}_c^*) \cdot \left(\frac{q\mathbf{E}}{T_0 s^2} - \nabla \ln B \right) + \xi^2 \frac{\mathbf{E} \times \mathbf{B}}{B^2} \cdot \nabla \ln B \right\} \\
 &\quad - \xi(1 - \xi^2) \left\{ \frac{\mu_0}{2B^2} \mathbf{J}_{\parallel} \cdot \mathbf{E} + \frac{T_0 s^2}{q} \mathbf{b} \cdot \nabla \left(\frac{\mu_0 J_{\parallel}}{B^2} \right) \right\} + (1 - \xi^2) \frac{T_0 s}{v_0 q B} \nabla \cdot \left(\mathbf{b} \times \frac{\partial \mathbf{b}}{\partial t} \right)
 \end{aligned}$$

This agrees exactly with Hazeltine's DKE, and includes additional terms in red. Green term enters from μ and is called the "twist" term by Ramos.

Linearization of δf DKE

Assuming $\mathbf{E}_0 = 0$ and using a speed norm v_0 that is constant in time yields

$$\begin{aligned}
 \delta \mathbf{v}_{gc} &= v_0 s \xi \mathbf{b}_1 + \frac{\mathbf{E}_1 \times \mathbf{B}_0}{B_0^2} + \frac{T_0 s^2}{q B_0^3} (1 + \xi^2) \left[(\mathbf{I} - 4\mathbf{b}_0 \mathbf{b}_0) \cdot \mathbf{B}_1 \times \nabla B_0 + \mathbf{b}_0 \times \nabla (\mathbf{B}_0 \cdot \mathbf{B}_1) \right] \\
 &+ \frac{T_0 s^2}{q B_0^2} \left\{ (1 - 3\xi^2) \delta(\mathbf{b}\mathbf{b}) \cdot \nabla \times \mathbf{B}_0 + \left[2\xi^2 (\mathbf{I} - \mathbf{b}_0 \mathbf{b}_0) + (1 - \xi^2) \mathbf{b}_0 \mathbf{b}_0 \right] \cdot \nabla \times \mathbf{B}_1 \right\} \\
 &- \frac{2T_0 s^2}{q} \frac{\mathbf{b}_0 \cdot \mathbf{B}_1}{B_0^3} \left[2\xi^2 (\mathbf{I} - \mathbf{b}_0 \mathbf{b}_0) + (1 - \xi^2) \mathbf{b}_0 \mathbf{b}_0 \right] \cdot \nabla \times \mathbf{B}_0 + \frac{mv_0 s \xi}{q B_0^2} \mathbf{b}_0 \times \frac{\partial \mathbf{B}_1}{\partial t} \\
 &= v_0 s \xi \mathbf{b}_1 + \frac{\mathbf{E}_1 \times \mathbf{B}_0}{B_0^2} + \frac{T_0 s^2}{q B_0^3} (1 + \xi^2) \left[(\mathbf{I} - 4\mathbf{b}_0 \mathbf{b}_0) \cdot \mathbf{B}_1 \times \nabla B_0 + \mathbf{b}_0 \times \nabla (\mathbf{B}_0 \cdot \mathbf{B}_1) \right] \\
 &+ \frac{2T_0 s^2}{q B_0^2} \left\{ \xi^2 \left(\mathbf{J}_{\perp 1} - \delta(\mathbf{b}\mathbf{b}) \cdot \mathbf{J}_0 - \frac{2\mathbf{B}_0 \cdot \mathbf{B}_1}{B_0^2} \mathbf{J}_{\perp 0} \right) + \frac{1 - \xi^2}{2} \left(\mathbf{J}_{\parallel 1} + \delta(\mathbf{b}\mathbf{b}) \cdot \mathbf{J}_0 - \frac{2\mathbf{B}_0 \cdot \mathbf{B}_1}{B_0^2} \mathbf{J}_{\parallel 0} \right) \right\} \\
 &+ \frac{mv_0 s \xi}{q B_0^2} \mathbf{b}_0 \times \frac{\partial \mathbf{B}_1}{\partial t}
 \end{aligned}$$

where $\mathbf{b}_1 = (\mathbf{I} - \mathbf{b}_0 \mathbf{b}_0) \cdot \mathbf{B}_1 / B_0$ and $\delta(\mathbf{b}\mathbf{b}) = \mathbf{b}_1 \mathbf{b}_0 + \mathbf{b}_0 \mathbf{b}_1$.

Linearization of δf DKE (cont.)

Linearizing the acceleration terms yields

$$\begin{aligned}
 \delta \dot{s} &= -\frac{s(1-\xi^2)}{2B_0^2} \mathbf{B}_0 \cdot \nabla \times \mathbf{E}_1 + \frac{q}{2sT_0} \mathbf{v}_{gc} \cdot \mathbf{E}_1 - s\delta \mathbf{v}_{gc} \cdot \nabla \ln v_0 \\
 &= -\frac{s(1-\xi^2)}{2B_0^2} \mathbf{B}_0 \cdot \nabla \times \mathbf{E}_1 + \frac{q}{2T_0s} (\mathbf{v}_{\parallel} + \mathbf{v}_c)_0 \cdot \mathbf{E}_1 + \frac{s}{2} (1+\xi^2) \frac{\mathbf{E}_1 \times \mathbf{B}_0}{B_0^2} \cdot \nabla \ln B_0 \\
 &\quad - s\delta \mathbf{v}_{gc} \cdot \nabla \ln v_0 \\
 \delta \dot{\xi} &= \frac{1-\xi^2}{2\xi} \left\{ \frac{\xi^2 \mathbf{B}_0}{B_0^2} \cdot \nabla \times \mathbf{E}_1 + (\mathbf{v}_{\parallel} + \mathbf{v}_c^*) \cdot \left(\frac{q\mathbf{E}_1}{T_0s^2} - \nabla \ln B_1 \right) - \delta (\mathbf{v}_{\parallel} + \mathbf{v}_c^*) \cdot \nabla \ln B_0 \right. \\
 &\quad \left. + \xi^2 \frac{\mathbf{E}_1 \times \mathbf{B}_0}{B_0^2} \cdot \nabla \ln B_0 - \frac{\xi^2 \mu_0}{B_0^2} \mathbf{J}_{\parallel 0} \cdot \mathbf{E}_1 - 2 \frac{T_0 s^2 \xi^2}{q} \left(\mathbf{b}_0 \cdot \nabla \delta \left(\frac{\mu_0 J_{\parallel}}{B^2} \right) + \mathbf{b}_1 \cdot \nabla \left(\frac{\mu_0 J_{\parallel 0}}{B_0^2} \right) \right) \right. \\
 &\quad \left. + \frac{2T_0 \xi s}{v_0 q B_0} \nabla \cdot \left(\mathbf{b}_0 \times \frac{\partial \mathbf{b}_1}{\partial t} \right) \right\}
 \end{aligned}$$

Continuum CEL-DKEs implemented in NIMROD for NTM studies

Qualities of Chapman-Enskog like (CEL) method*:

- ▶ Separates fluid and kinetic parts of distribution function
- ▶ Fluid equations govern lowest order fluid quantities, n_a , \mathbf{V}_a , and T_a
- ▶ Kinetic equation governs kinetic distortion, F_a
- ▶ n_a , \mathbf{V}_a , and T_a provide thermodynamic drives for F_a
- ▶ Moments of F_a close fluid equations

Research Objective: Understand challenges

- ▶ Strong nonlinear coupling between fluid and F_a
- ▶ Scaling velocity by thermal speed
- ▶ Implicit advance for large time steps

*S. Chapman and T.G. Cowling, The Mathematical Theory of Non-Uniform Gases (Cambridge University Press, Cambridge, 1939); Z. Chang and J.D. Callen, Phys. Fluids 4, 1167 (1992).

CEL method separates fluid and kinetic physics

Starting from the DKE* project out Maxwellian part, $f = f^M + F$,
and transform to coordinates, $(s, \xi) \equiv (|\mathbf{v} - \mathbf{V}|/v_T, \mathbf{v} \cdot \mathbf{B}/|\mathbf{v}||\mathbf{B}|)$:

$$\frac{\partial F}{\partial t} + \mathbf{v}_{gc} \cdot \nabla F + \dot{s} \frac{\partial F}{\partial s} + \dot{\xi} \frac{\partial F}{\partial \xi} = C - f^M \left[\frac{d \ln n}{dt} + \frac{2s}{v_T} \cdot \frac{d\mathbf{V}}{dt} + \left(s^2 - \frac{3}{2} \right) \frac{d \ln T}{dt} \right]$$

where

$$\begin{aligned} \mathbf{v}_{gc} &= v_T s \xi \mathbf{b} + \frac{\mathbf{E} \times \mathbf{B}}{B^2} + \frac{T s^2}{q B} (1 + \xi^2) \mathbf{b} \times \nabla \ln B \\ &\quad + \frac{2T s^2}{q B^2} \left[\xi^2 (\mathbf{I} - \mathbf{b}\mathbf{b}) + \frac{1}{2} (1 - \xi^2) \mathbf{b}\mathbf{b} \right] \cdot \nabla \times \mathbf{B} + \frac{m v_T s \xi}{q B^2} \mathbf{b} \times \frac{\partial \mathbf{B}}{\partial t} \\ \dot{s} &= -\frac{s}{2} \frac{d \ln T}{dt} + \frac{s(1 - \xi^2)}{2} \frac{\partial \ln B}{\partial t} + \frac{q \mathbf{v}_{gc} \cdot \mathbf{E}}{2sT} \\ \dot{\xi} &= \frac{1 - \xi^2}{2\xi} \left\{ -\xi^2 \frac{\partial \ln B}{\partial t} + (\mathbf{v}_{\parallel} + \mathbf{v}_c^*) \cdot \left(\frac{q\mathbf{E}}{T s^2} - \nabla \ln B \right) + \xi^2 \frac{\mathbf{E} \times \mathbf{B}}{B^2} \cdot \nabla \ln B \right. \\ &\quad \left. - \frac{\xi^2}{B^2} [\mathbf{b}\mathbf{b} \cdot (\nabla \times \mathbf{B})] \cdot \mathbf{E} - 2 \frac{T s^2 \xi^2}{q} \mathbf{b} \cdot \nabla \left[\frac{\mathbf{b} \cdot (\nabla \times \mathbf{B})}{B^2} \right] + \frac{m v_T s \xi}{q B} \nabla \cdot \left(\mathbf{b} \times \frac{\partial \mathbf{b}}{\partial t} \right) \right\} \\ \mathbf{v}_c^* &= \frac{2T s^2 \xi^2}{q B^2} (\mathbf{I} - \mathbf{b}\mathbf{b}) \cdot \nabla \times \mathbf{B} + \frac{m v_T s \xi}{q B^2} \mathbf{b} \times \frac{\partial \mathbf{B}}{\partial t} \end{aligned}$$

*R.D. Hazeltine, Plasma Phys. 15, 77 (1973); R.D. Hazeltine and J.D. Meiss, Plasma Confinement (Adison-Wesley, Redwood City, 1992); J. J. Ramos, Phys Plasmas 17, 082502 (2010).

Challenges highlighted in kinetic thermal transport case studies

$$\frac{3}{2}n \frac{\partial T}{\partial t} = \kappa_{\perp} \nabla \cdot [(\mathbf{I} - \mathbf{b}\mathbf{b}) \cdot \nabla T] - \nabla \cdot \mathbf{q}_{\parallel} + Q$$

Calculate parallel heat flux as moment of kinetic distortion

$$\mathbf{q}_{\parallel} = \frac{m}{2} \int d\mathbf{v} v^2 v_{\parallel} F = \pi m v_T^6 \int_{-1}^1 d\xi \int_0^{\infty} ds (s^5 \xi F)$$

$$\begin{aligned} & \frac{\partial F}{\partial t} + \mathbf{v}_{\parallel} \cdot \nabla F - \frac{1 - \xi^2}{2\xi} \mathbf{v}_{\parallel} \cdot \nabla \ln B \frac{\partial F}{\partial \xi} - \frac{s}{2} \left(\mathbf{v}_{\parallel} \cdot \nabla + \frac{\partial}{\partial t} \right) \ln T \frac{\partial F}{\partial s} \\ & = C + \left(\frac{5}{2} - s^2 \right) \mathbf{v}_{\parallel} \cdot \nabla \ln T f^M + \frac{2}{3nT} \left(s^2 - \frac{3}{2} \right) (\nabla \cdot \mathbf{q}_{\parallel} - Q) f^M \end{aligned}$$

(red terms have temperature dependence.)

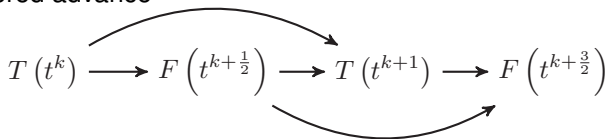
Possible θ -centered semi-implicit time advances

Problem: tight nonlinear coupling of fluid and kinetic distortion

$$\begin{aligned}\frac{\partial T}{\partial t} &= G(T, F) \\ \frac{\partial F}{\partial t} &= H(T, F)\end{aligned}$$

complex nonlinear combinations of T and F

► Staggered advance



$$\Delta T - \theta \Delta t G_{\text{lin}}(\Delta T, F^{k+\frac{1}{2}}) = \Delta t G(T^k, F^{k+\frac{1}{2}})$$

$$\Delta F - \theta \Delta t H_{\text{lin}}(T^{k+1}, \Delta F) = \Delta t H(T^{k+1}, F^{k+\frac{1}{2}})$$

► Simultaneous advance (**Picard iterations** or **Newton iterations**)

$$T(t^k), F(t^k) \longrightarrow T(t^{k+1}), F(t^{k+1})$$

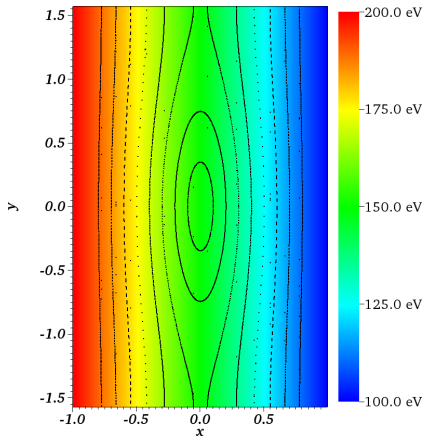
$$\begin{aligned}\Delta T - \theta \Delta t G(T^{k+1}, F^{k+1}) &= (1 - \theta) \Delta t G(T^k, F^k) \\ \Delta F - \theta \Delta t H(T^{k+1}, F^{k+1}) &= (1 - \theta) \Delta t H(T^k, F^k)\end{aligned}$$

GMRES fails to solve

Test case: thermal transport in magnetic island

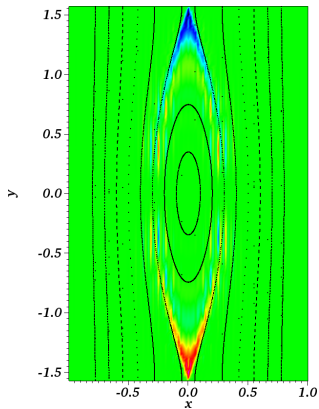
Kinetic parallel thermal transport across magnetic island in slab geometry

- ▶ $n = 9.5175 \times 10^{18} \text{ m}^{-3}$, $\mathbf{V} = 0$
- ▶ Ignore electron-ion and ion-electron collisions
- ▶ Boundary condition: periodic in Z direction
- ▶ Objective: take as large time steps as possible to get to steady state with kinetic parallel heat flux
- ▶ 32x32 grid in xy -plane
- ▶ 3rd degree polynomials

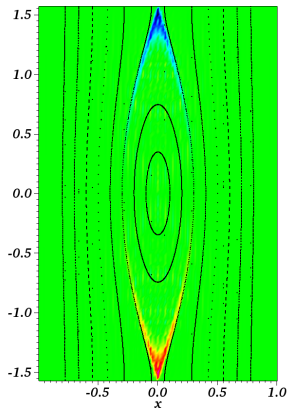
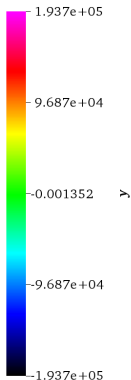


Initial temperature is a linear gradient that flattens across island as T evolves

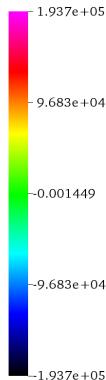
Standard and mixed finite element steady state parallel heat flux with conductivity $\kappa_{\parallel} = 1.5 \times 10^7$



Standard fluid steady state
 q_{\parallel} [W/m^2]



Mixed finite element steady state
 q_{\parallel} [W/m^2]



Picard and Newton methods implemented for simultaneous T and F advance.

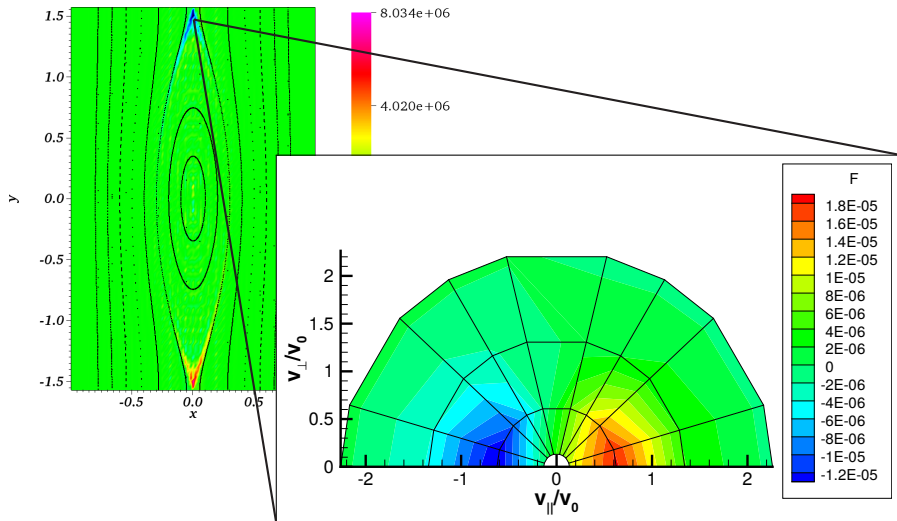
Implicit advance of F:

$$\begin{aligned} & \frac{F^{k+1} - F^k}{\Delta t} + \sqrt{\frac{2T}{m}} s\xi \left(\nabla_{\parallel} F^{k+1} - \frac{1 - \xi^2}{2\xi} \nabla_{\parallel} \ln B \frac{\partial F^{k+1}}{\partial \xi} \right) - \frac{s}{2T} \left(\sqrt{\frac{2T}{m}} s\xi \nabla_{\parallel} + \frac{\partial}{\partial t} \right) T \frac{\partial F^{k+1}}{\partial s} \\ & = C(T, F^{k+1}) + \left[\left(\frac{5}{2} - s^2 \right) \sqrt{\frac{2}{mT}} s\xi \nabla_{\parallel} T + \frac{2}{3nT} \left(s^2 - \frac{3}{2} \right) (\nabla \cdot \mathbf{q}_{\parallel} (F^{k+1}, T) - G) \right] f^M(T) \end{aligned}$$

Implicit advance of T:

$$\frac{3}{2} n \frac{T^{k+1} - T^k}{\Delta t} = \kappa_{\perp} \nabla \cdot \left[(\mathbf{I} - \mathbf{b}\mathbf{b}) \cdot \nabla T^{k+1} \right] - \nabla \cdot \mathbf{q}_{\parallel} (T, F) + G(T, F)$$

Kinetic heat flux calculated as moment of distribution function



Newton more costly than Picard iterations but can take larger time step

- ▶ 256 processors, 32x32 grid, polynomial degree=3
- ▶ Starting from MFE steady state, run an additional 10^{-5} s

	Δt	wall clock time to $t = 10^{-5}$ s	average GMRES iterations per step	time per iteration
Picard	10^{-8} s	75 mins	5	0.9 s
Newton	10^{-8} s	200 mins	4	3 s
Newton	10^{-7} s	49 mins	52	0.57 s
Newton	10^{-6} s	42 mins	723	0.35 s

- ▶ Need to implement parallelism over speed grid points for efficiency improvement.

Upcoming work

- ▶ Implement s-parallelism for simultaneous advance of T and F .
- ▶ Possibly speed-up Newton iterations
(reuse preconditioning matrix, improved check for convergence)
- ▶ Adaptive time step
- ▶ Examine needed velocity grid for electron-ion collisions
- ▶ Use developed code in a tearing mode simulation
with evolving \mathbf{B} , n , \mathbf{V} .